

Preequilibrium Cluster Emission and Systematics on the Formation Cross Sections of Superheavy Nuclei (SHN)

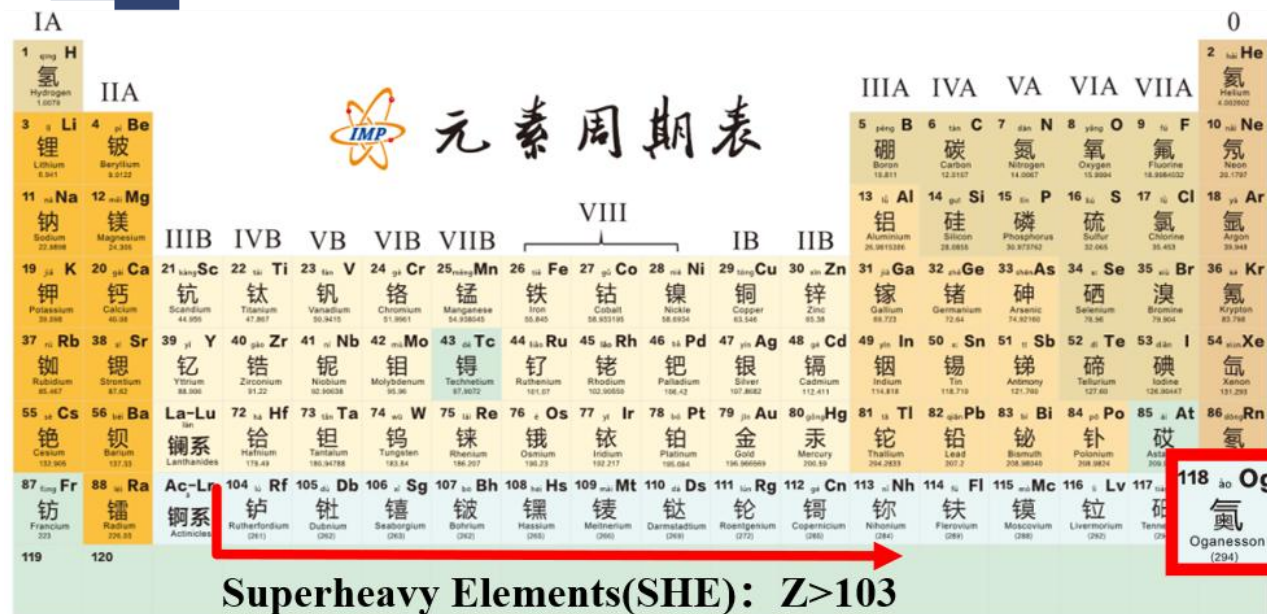
Zi-Han Wang (王子涵)

Supervisor: Zhao-Qing Feng (冯兆庆)

Outline

- 1 Research Motivation
- 2 Preequilibrium Cluster Emission
- 3 Formation Cross Sections of SHN

1. Research Motivation: superheavy nuclei

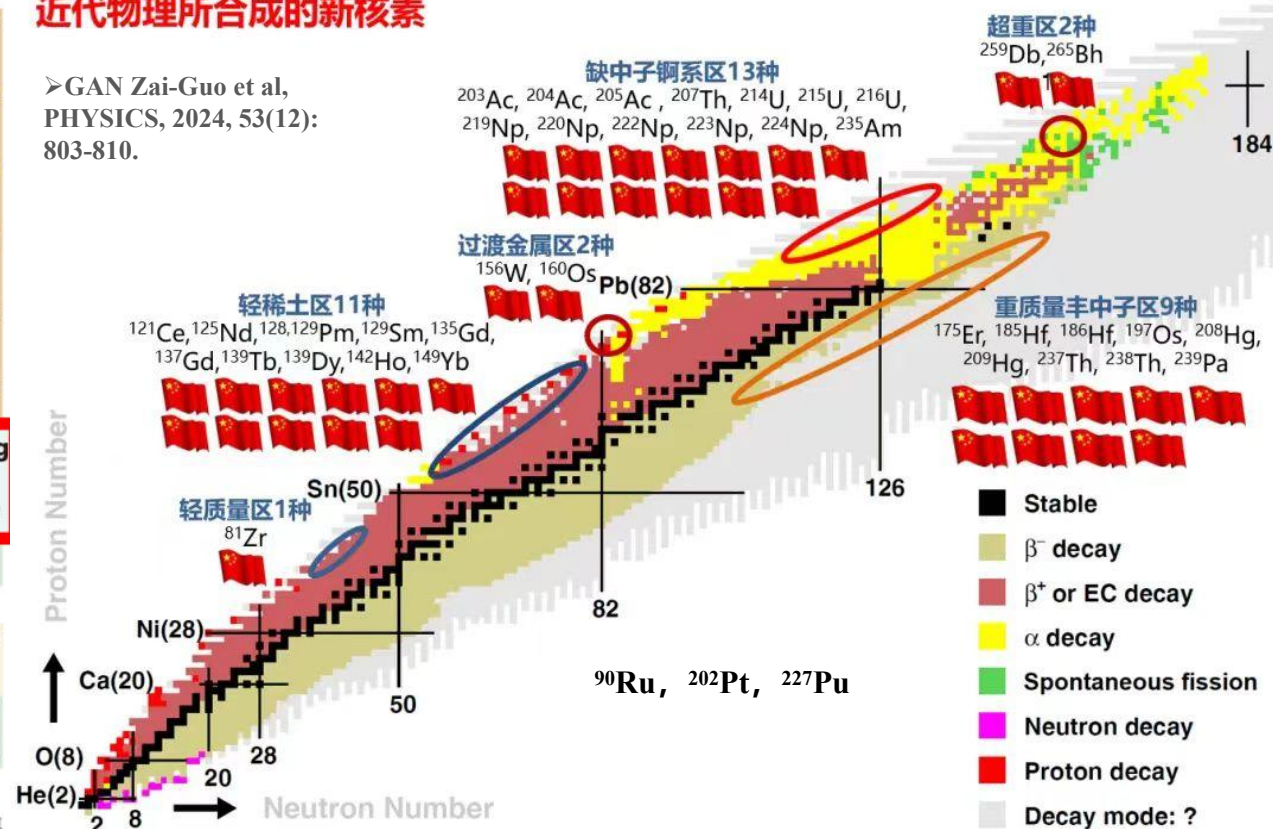


Significance of SHN research:

- Expand the periodic table of elements
- Explore the limits of atomic nucleus mass and charge
- Study atomic nucleus decay and structural properties
- Test the shell model theory of atomic nuclei
- Reveal the nuclear synthesis process in nuclear astrophysics

近代物理所合成的新核素

➢ GAN Zai-Guo et al,
PHYSICS, 2024, 53(12):
803-810.

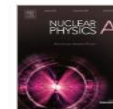


Difficulties in the synthesis of new SHN:

- Extremely low production cross-section and fleeting lifetime
- The products are difficult to separate and identify
- The experimental cost is high and the cycle is long

Accurate Theoretical Predictions Are Needed !

1. Research Motivation: MNT reactions



Nuclear Physics A
Volume 176, Issue 2, 29 November 1971, Pages 284-288

In the multi-nucleon transfer (MNT) reactions:

- Generate broad isotope distribution depending on transfer channels.
- Angular and excitation energy of compound nucleus distribute widely.
- Producing neutron-rich heavy and superheavy nuclei.

弹核碎裂/深部非弹反应 Projectile fragmentation/Deep inelastic

质谱/原子光谱/阴极射线技术 Mass spectroscopy/AS/CR

放射性衰变 Radioactive decay

融合/转移反应 Fusion/Transfer

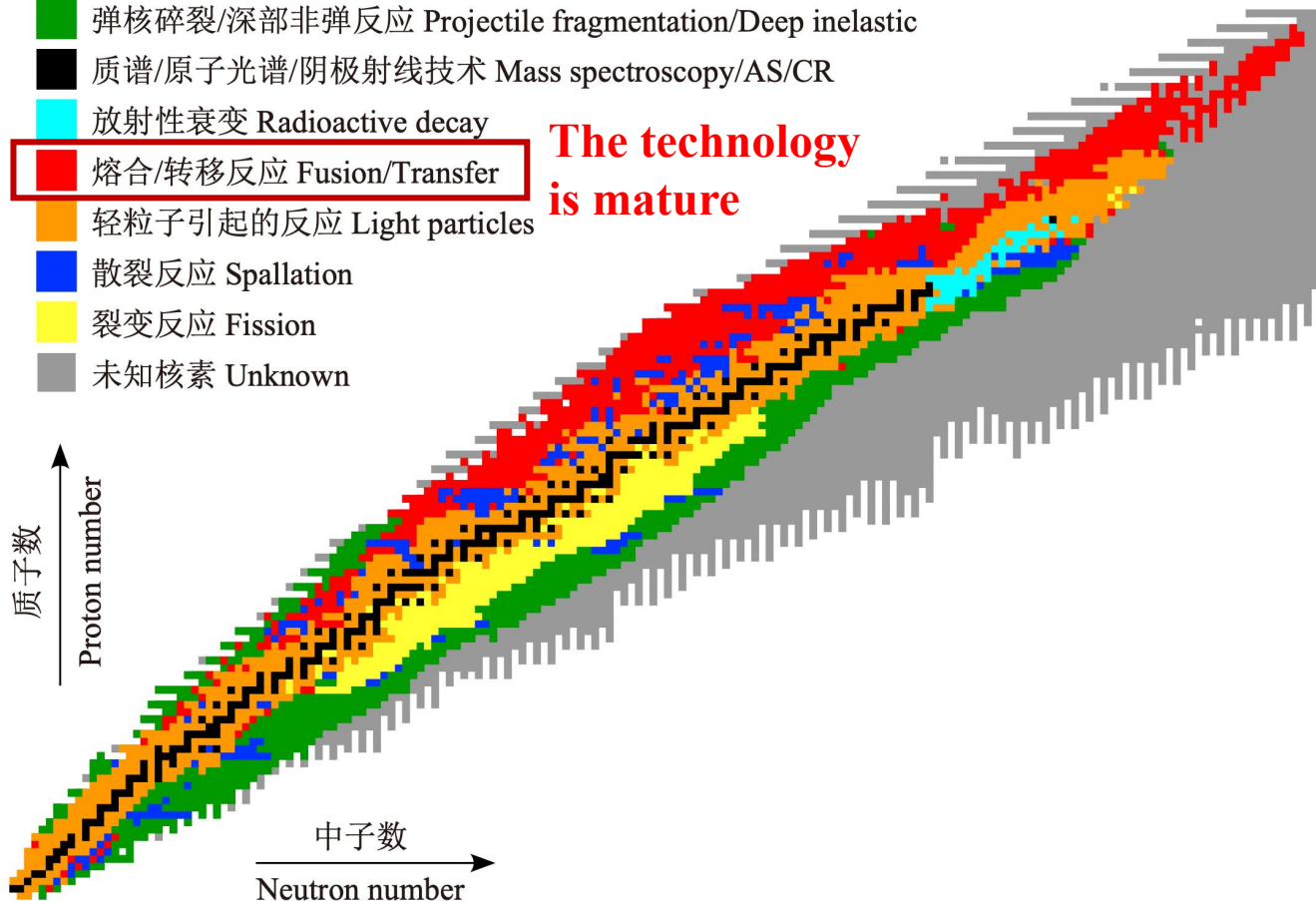
轻粒子引起的反应 Light particles

散裂反应 Spallation

裂变反应 Fission

未知核素 Unknown

The technology
is mature



New isotopes $^{29,30}\text{Mg}$, $^{31,32,33}\text{Al}$, $^{33,34,35,36}\text{Si}$, $^{35,36,37,38}\text{P}$, $^{39,40}\text{S}$ and $^{41,42}\text{Cl}$ produced in bombardment of a ^{232}Th target with 290 MeV ^{40}Ar ions

A.G. Artukh, V.V. Avdeichikov [†], G.F. Gridnev, V.L. Mikheev, V.V. Volkov, J. Wilczyński

Multi-nucleon transfer reactions (MNT) experiment progress

Lab.	Reaction system	References
Dubna	$^{136}\text{Xe} + ^{208}\text{Pb}$	Phys. Rev. C 86, 044611 (2012)
	$^{156,160}\text{Gd} + ^{186}\text{W}$	Phys. Rev. C 96, 064621 (2017)
GSI	$^{238}\text{U} + ^{238}\text{U}$	Phys. Rev. Lett. 39, 385 (1977)
	$^{48}\text{Ca} + ^{248}\text{Cm}$	Phys. Rev. Lett. 41, 469 (1978)
	$^{48}\text{Ca} + ^{238}\text{U}$	Phys. Lett. B 748, 199 (2015)
		Eur. Phys. J. A 56, 224 (2020)
GANIL	$^{238}\text{U} + ^{238}\text{U}$	IJMPE 17, 2235-2239 (2008)
	$^{136}\text{Xe} + ^{198}\text{Pt}$	Phys. Rev. Lett. 115, 172503 (2015)
Argonne	$^{136}\text{Xe} + ^{208}\text{Pb}$	Phys. Rev. C 91, 064615 (2015)
	$^{204}\text{Hg} + ^{198}\text{Pt}$	Physics Letters B 771, 119-124 (2017)
RIKEN	$^{238}\text{U} + ^{198}\text{Pt} \rightarrow ^{241}\text{U}$	Phys. Rev. Lett. 130, 132502 (2023)



1. Research Motivation: pre-equilibrium clusters



□ Theoretically

➤ The Exciton Model

- M. Blann, Rev. Nucl. Sci. 25, 123 (1975).
- E. Gadioli et al, Phys. Rev. C 16, 1404 (1977).
- OV Fotina et al, Int. J. Mod. Phys. E 19, 1134 (2010).

➤ The Dinuclear System (DNS) Model

- Z. Q. Feng, Phys. Rev. C 107, 054613 (2023).

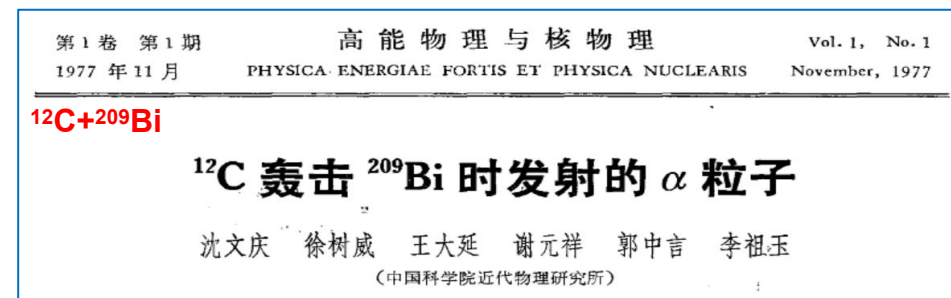
➤ The significance of studying pre-equilibrium clusters

- ✓ the correlation of spatial configuration of nucleons
- ✓ the cluster structure of the stable or unstable nuclide
- ✓ the cluster formation mechanism in nuclear reactions
- ✓ the kinetic mechanism of MNT reactions
- ✓

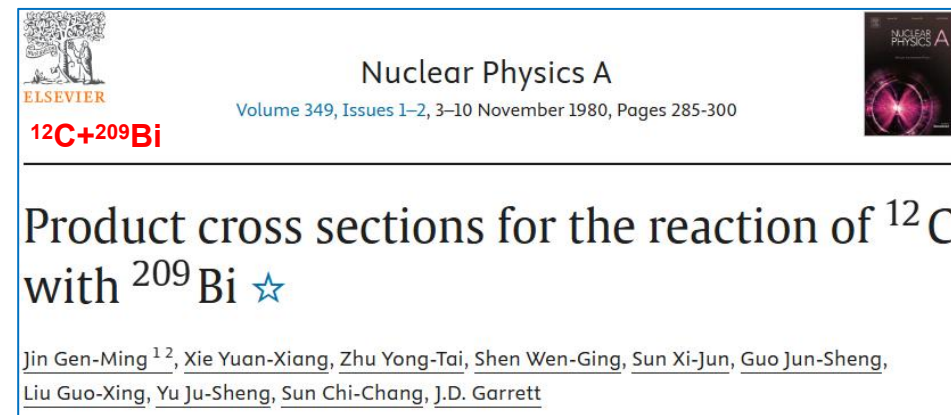
- P. E. Hodgson and E. Běták, Phys. Rep. 374, 1-89 (2003).

□ Experimentally

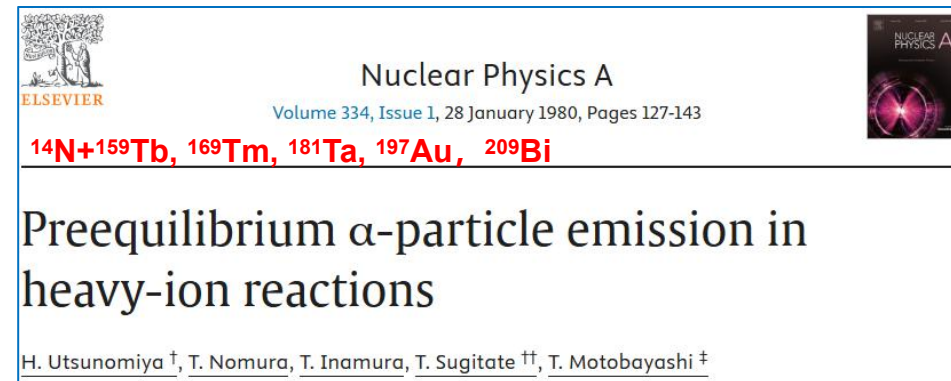
➤ 1977, IMP:



➤ 1980, IMP:



➤ 1980, RIKEN:

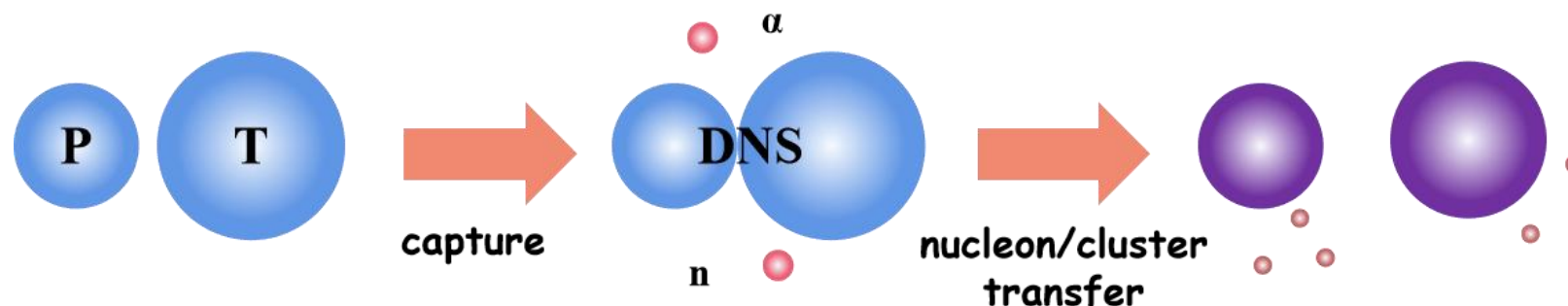


2. Preequilibrium Cluster Emission

arXiv:2505.02337



□ The Dinuclear System Model (DNS)



Schematic diagram of multinucleon transfer reactions in the DNS model

➤ 冯兆庆. 近垒重离子熔合反应和超重核合成机制研究[D]. 甘肃:中国科学院近代物理研究所,2007.

➤ Emission cross-section of pre-equilibrium cluster:

$$\sigma_v(E_k, \theta, t) = \sum_{J=0}^{J_{\max}} \sum_{Z_1=Z_v}^{Z_{\max}} \sum_{N_1=N_v}^{N_{\max}} \sigma_{cap}(E_{c.m.}, J) \times \int f(B) \times P(Z_1, N_1, E_1(E_{c.m.}, J), t, \beta, B) \times P_v(Z_v, N_v, E_k) dB$$

Emission probability
of the cluster:

$$P_v(Z_v, N_v, E_k) = \Delta t \frac{\Gamma_v}{\hbar}$$

$$v = n, p, d, t, {}^3\text{He}, \alpha, {}^6, {}^7\text{Li}, {}^8, {}^9\text{Be}$$

2. Preequilibrium Cluster Emission

arXiv:2505.02337



□ The Dinuclear System Model (DNS)—Capture process:

Capture cross section:

$$\sigma_{cap}(E_{c.m.}, J) = \frac{\pi \hbar^2 (2J+1) T(E_{c.m.}, J)}{2 \mu E_{c.m.}} \quad \mu = m_n \frac{A_P A_T}{A_P + A_T}$$

Penetration probability:

$$T(E_{c.m.}, J) = \int f(B) T(E_{c.m.}, J, B) dB$$

1. for the light and medium systems: the Hill-Wheeler formula

$$T(E_{c.m.}, J, B) = \left\{ 1 + \exp \left\{ -\frac{2\pi}{\hbar \omega(J)} \left[E_{c.m.} - B - \frac{\hbar^2 J(J+1)}{2 \mu R_B^2(J)} \right] \right\} \right\}^{-1}$$

➤ D. L. Hill and J. A. Wheeler, Phys. Rep. 89, 1102-1145 (1953).

Barrier distribution function:

Gaussian form $f(B) = \frac{1}{N} \exp \left(-\left(\frac{B - B_m}{\Delta} \right)^2 \right)$

$$\Delta = (B_C - B_S) / 2 \quad B_m = (B_C + B_S) / 2 \quad \int f(B) dB = 1$$

2. for the heavy systems: the classical trajectory approach

$$T(E_{c.m.}, J, B) = 0, \quad E_{c.m.} < B + \frac{J(J+1)\hbar^2}{2\mu R_C^2}; \quad T(E_{c.m.}, J, B) = 1, \quad E_{c.m.} > B + \frac{J(J+1)\hbar^2}{2\mu R_C^2}; \quad R_C = r_{0c} (A_P^{1/3} + A_T^{1/3})$$

2. Preequilibrium Cluster Emission

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□ The Dinuclear System Model (DNS)—Transfer dynamics:

Master equation:

Nucleons (**n**, **p**) + Clusters (**d**, **t**, **³He**, **⁴He**):

Only the transfer of Nucleons:

$$\begin{aligned} \frac{dP(Z_1, N_1, E_1, t)}{dt} &= \sum_{Z'_1} W_{Z_1, N_1; Z'_1, N_1}(t) [d_{Z_1, N_1} P(Z'_1, N_1, E'_1, t) \\ &\quad - d_{Z'_1, N_1} P(Z_1, N_1, E_1, t)] + \sum_{N'_1} W_{Z_1, N_1; Z_1, N'_1}(t) \\ &\quad \times [d_{Z_1, N_1} P(Z_1, N'_1, E'_1, t) - d_{Z_1, N'_1} P(Z_1, N_1, E_1, t)]. \end{aligned}$$



$$\begin{aligned} \frac{dP(Z_1, N_1, E_1, \beta_1, B, t)}{dt} &= \sum_{Z'_1=Z_1 \pm 1} W_{Z_1, N_1, \beta_1; Z'_1, N_1, \beta'_1}(t) [d_{Z_1, N_1} P(Z'_1, N_1, E'_1, \beta'_1, B, t) - d_{Z'_1, N_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \\ &\quad + \sum_{N'_1=N_1 \pm 1} W_{Z_1, N_1, \beta_1; Z_1, N'_1, \beta'_1}(t) [d_{Z_1, N_1} P(Z_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \\ &\quad + \sum_{Z'_1=\pm 1, N'_1=N_1 \pm 1} W_{Z_1, N_1, \beta_1; Z'_1, N'_1, \beta'_1}^d(t) [d_{Z_1, N_1} P(Z'_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z'_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \\ &\quad + \sum_{Z'_1=\pm 1, N'_1=N_1 \pm 2} W_{Z_1, N_1, \beta_1; Z'_1, N'_1, \beta'_1}^t(t) [d_{Z_1, N_1} P(Z'_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z'_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \\ &\quad + \sum_{Z'_1=\pm 2, N'_1=N_1 \pm 1} W_{Z_1, N_1, \beta_1; Z'_1, N'_1, \beta'_1}^{^3He}(t) [d_{Z_1, N_1} P(Z'_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z'_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \\ &\quad + \sum_{Z'_1=\pm 2, N'_1=N_1 \pm 2} W_{Z_1, N_1, \beta_1; Z'_1, N'_1, \beta'_1}^{\alpha}(t) [d_{Z_1, N_1} P(Z'_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z'_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)]. \end{aligned}$$

$$W_{Z_1, N_1; Z'_1, N_1} = \frac{\tau_{mem}(Z_1, N_1, E_1; Z'_1, N_1, E'_1)}{d_{Z_1, N_1} d_{Z'_1, N_1} \hbar^2} \sum_{i, i'} |\langle i' | V | i \rangle|^2$$

➤ Z. Q. Feng, Phys. Rev. C 108, L051601 (2023).

$$W_{Z_1, N_1; Z'_1, N_1}^s = G_s \frac{\tau_{mem}(Z_1, N_1, E_1; Z'_1, N_1, E'_1)}{d_{Z_1, N_1} d_{Z'_1, N_1} \hbar^2} \sum_{ii'} |\langle i' | V | i \rangle|^2$$

➤ R. Mattiello et al, Phys. Rev. C 55, 1443 (1997).
➤ Z. Q. Feng, Phys. Rev. C 102, 044604 (2020).

2. Preequilibrium Cluster Emission

arXiv:2505.02337



□ The Dinuclear System Model (DNS)—Preequilibrium cluster emission:

Emission probability : $P_\nu(Z_\nu, N_\nu, E_k) = \Delta t \frac{\Gamma_\nu}{\hbar}$

➤ the time step: $\Delta t = 0.5(0.25) \times 10^{-22} s$

➤ the particle decay widths Γ_ν :

the Weisskopf evaporation theory

$$\Gamma_\nu(E^*, J) = (2s_\nu + 1) \frac{m_\nu}{\pi^2 \hbar^2 \rho(E^*, J)} \times \int_0^{E^* - B_\nu - E_{rot} - V_c} \varepsilon \times \rho(E^* - B_\nu - E_{rot} - V_c - \varepsilon, J) \sigma_{inv}(\varepsilon) d\varepsilon$$

- V. Weisskopf, Phys. Rev. 52, 295 (1937).
- P. H. Chen et al, Chin. Phys. C 40, 091002 (2016).

Kinetic energy spectra:

the Monte Carlo method, $\varepsilon_\nu = (0, E^* - B_\nu - V_c - E_{rot})$

➤ for neutron: Watt Spectrum

$$\frac{dN_n}{d\epsilon_n} = C_n \frac{\epsilon_n^{1/2}}{T_w^{3/2}} \exp\left(-\frac{\epsilon_n}{T_w}\right) \quad T_w = 1.7 \pm 0.1 MeV$$

➤ H. Rossner, Phys. Rev. C 45, 719 (1992).

➤ for charged particles: Boltzmann Distribution

$$\frac{dN_\nu}{d\epsilon_\nu} = 8\pi E_k \left(\frac{m}{2\pi T_\nu}\right)^{1/2} \exp\left(-\frac{\epsilon_\nu}{T_\nu}\right) \quad T_\nu = \sqrt{\frac{E^*}{a}}, \quad a = A/12$$

2. Preequilibrium Cluster Emission

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□ The Dinuclear System Model (DNS)—Preequilibrium cluster emission:

Angular distributions : the deflection function method

$$\Theta(J_i) = \Theta_C(J_i) + \Theta_N(J_i)$$

➤ C. Peng and Z. Q. Feng, Eur. Phys. J. A 58, 162 (2022).

➤ The coulomb deflection

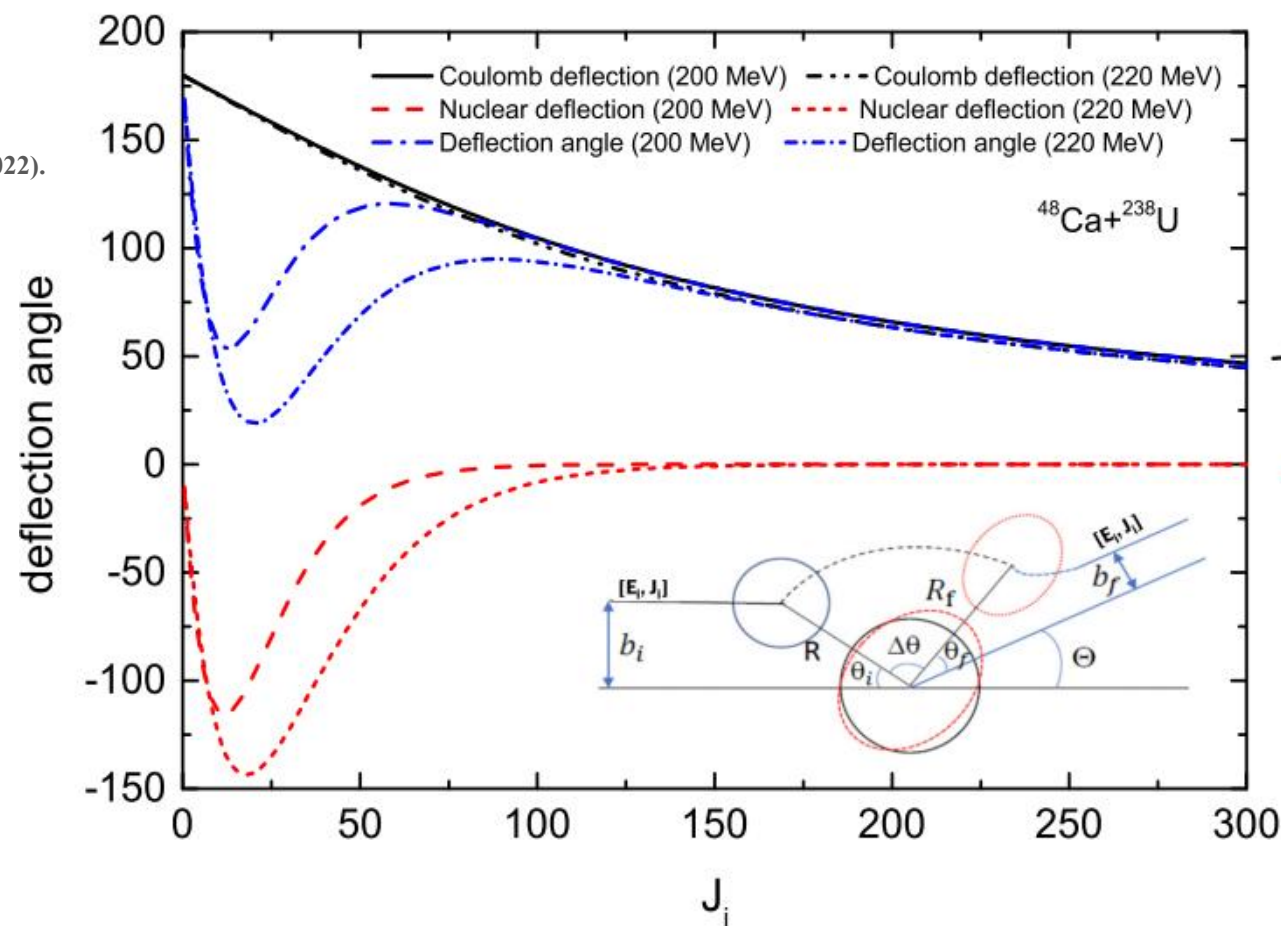
$$\Theta_C(J_i) = 2 \tan^{-1} \frac{Z_P Z_T e^2}{2 E_{c.m.} b}$$

➤ The nuclear deflection

$$\Theta_N(J_i) = -\beta \Theta_C^{gr}(J_i) \frac{J_i}{J_{gr}} \left(\frac{\delta}{\beta} \right)^{\frac{J_i}{J_{gr}}}$$

the grazing angular momentum :

$$J_{gr} = 0.22 R_{int} \left[A_{red} (E_{c.m.} - V(R_{int}))^{1/2} \right]$$



➤ Z. Q. Feng, Phys. Rev. C 107, 054613 (2023). 10

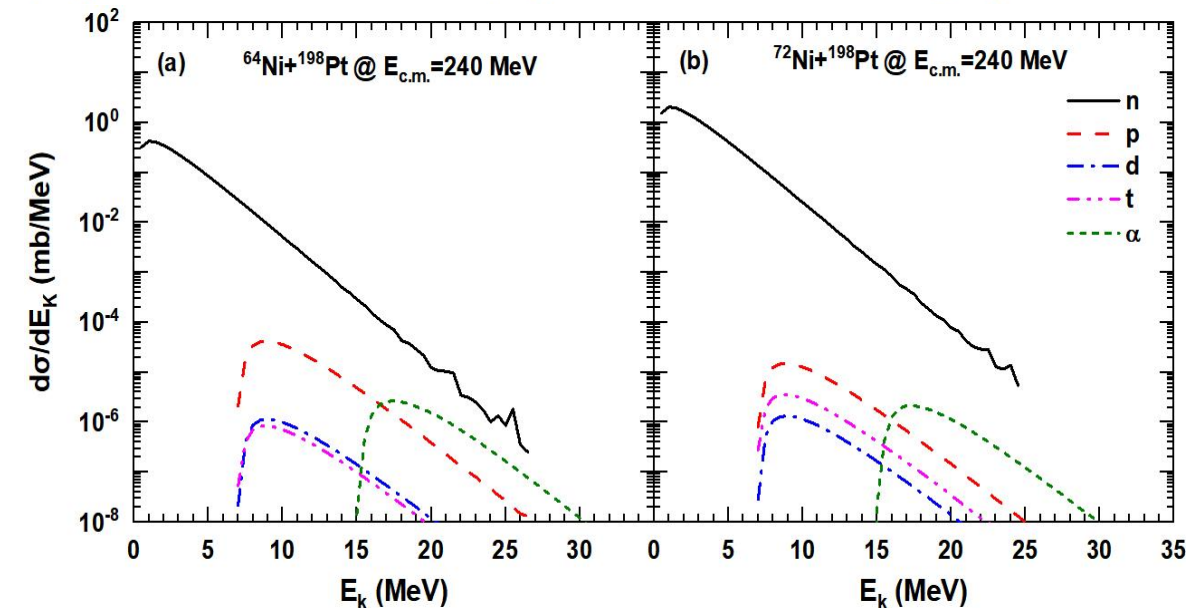
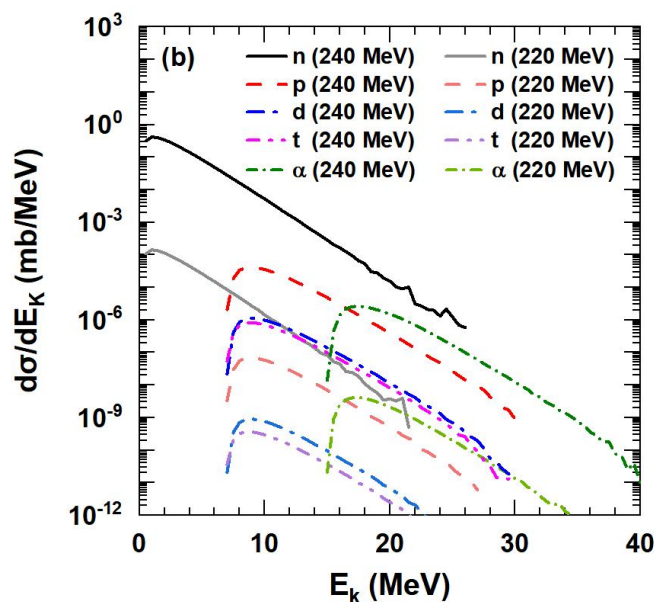
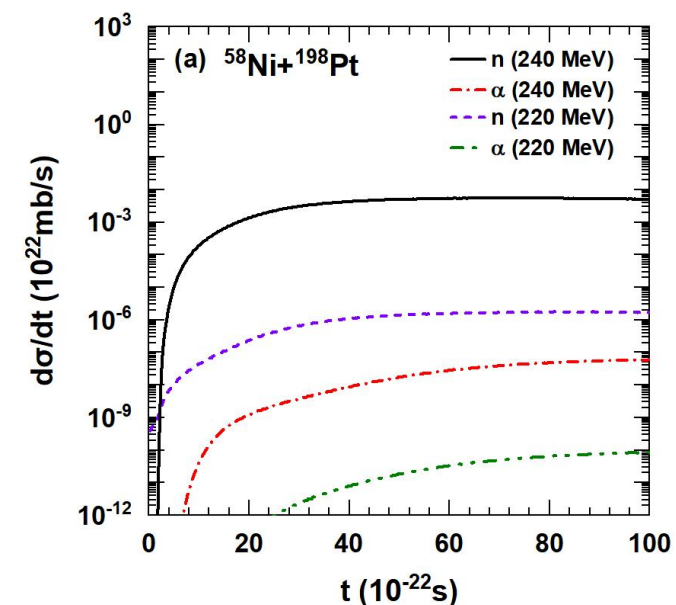
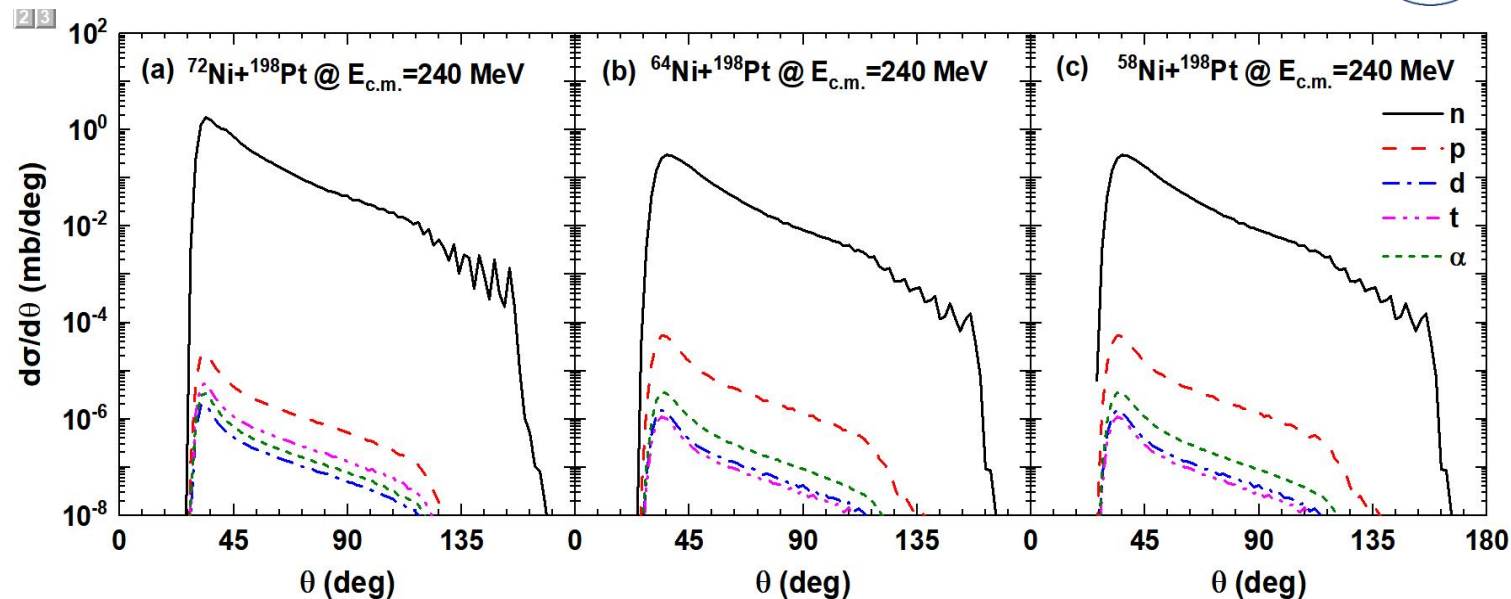
2. Preequilibrium Cluster Emission

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□ Results: $^{58,64,72}\text{Ni} + ^{198}\text{Pt}$

Angular distributions



Temporal evolution

Kinetic energy spectra

2. Preequilibrium Cluster Emission

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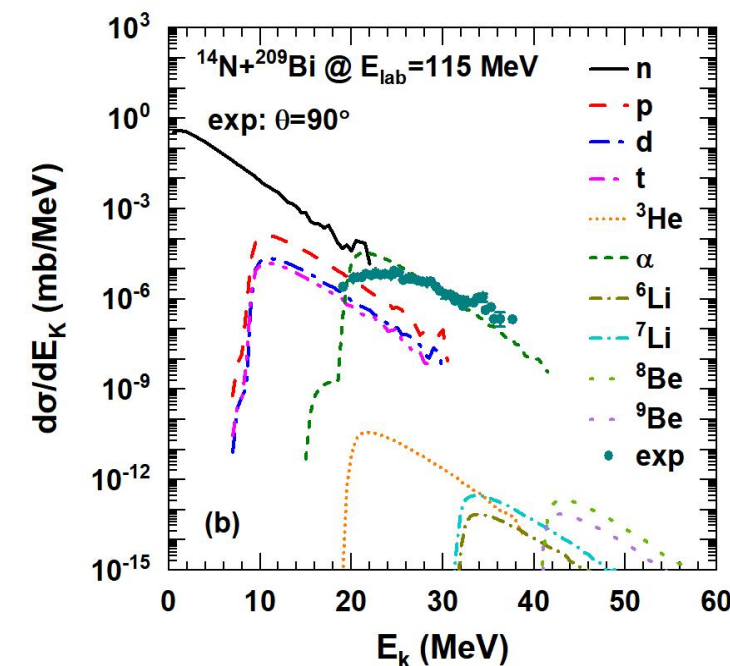
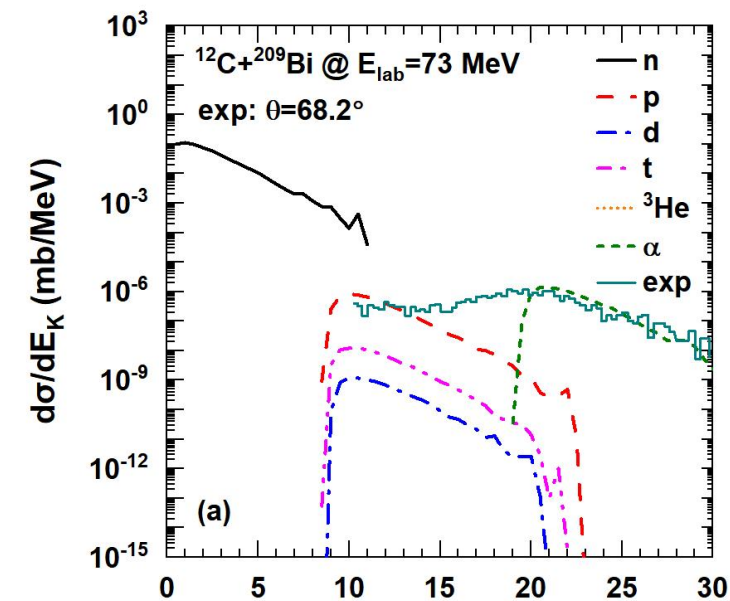
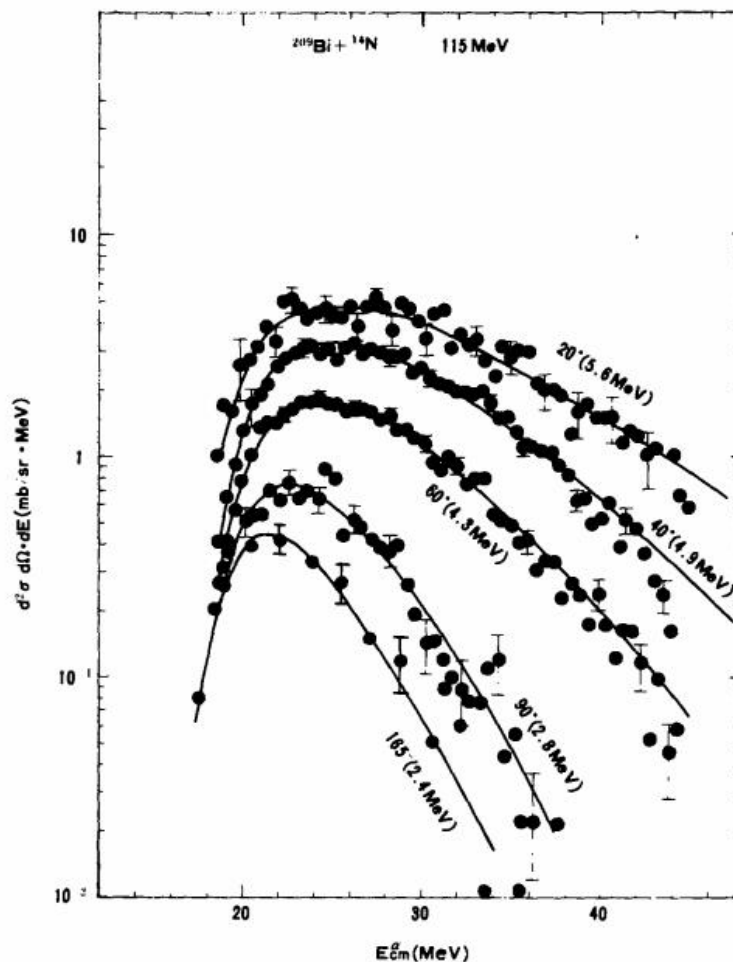
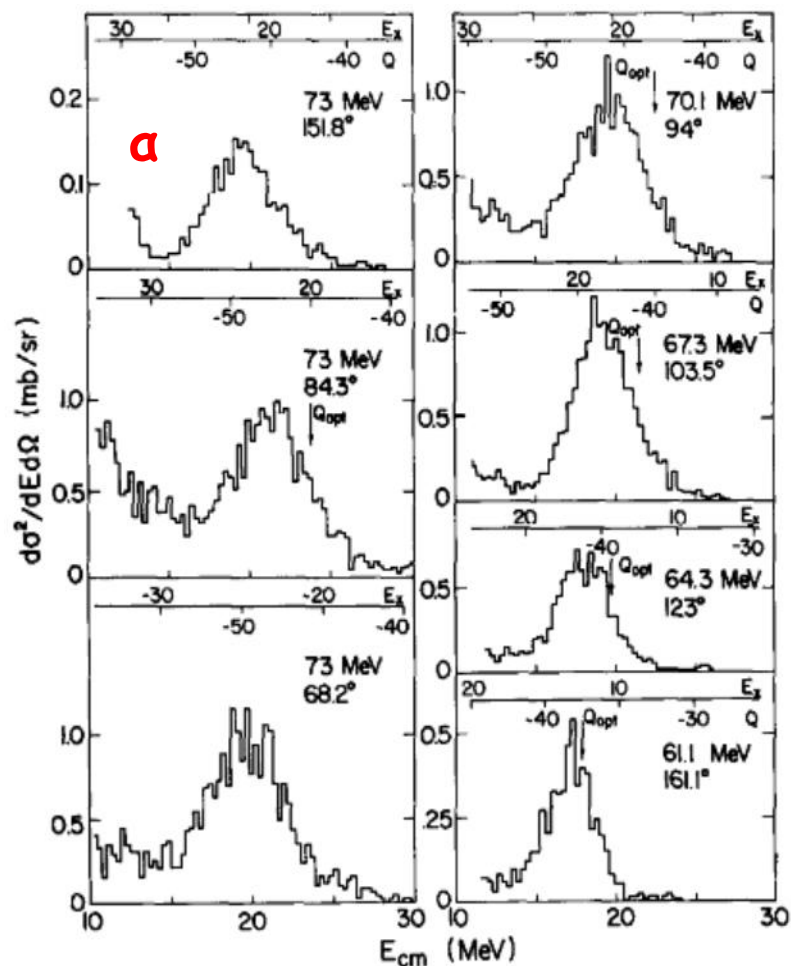


Results: ^{12}C , $^{14}\text{N}+^{209}\text{Bi}$

IMP: $^{12}\text{C}+^{209}\text{Bi}@E_{\text{lab}}=73\text{MeV}$

RIKEN:

$^{14}\text{N}+^{209}\text{Bi}@E_{\text{lab}}=115\text{MeV}$

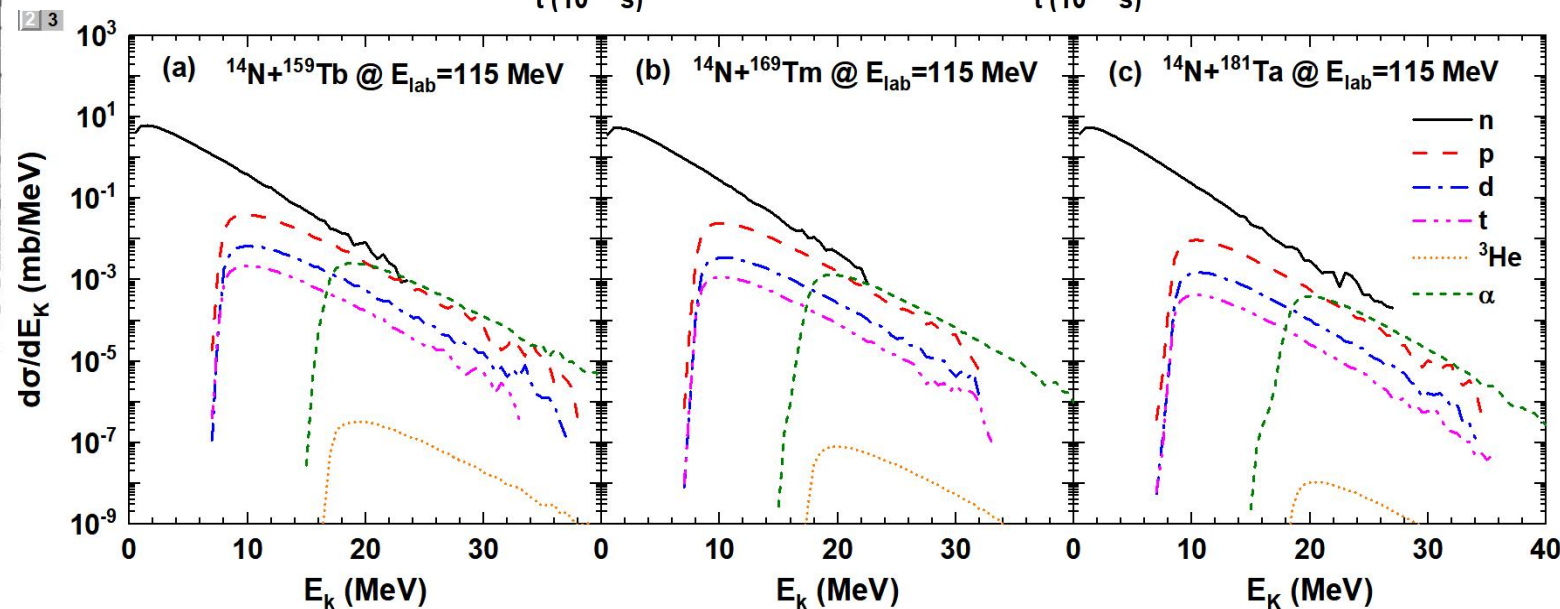
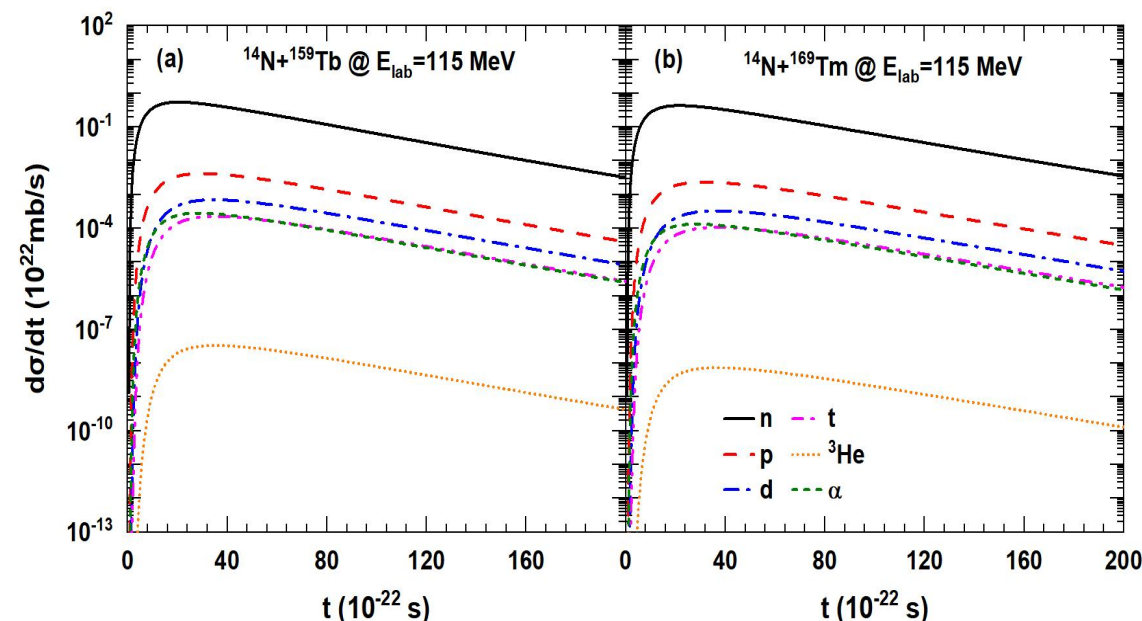
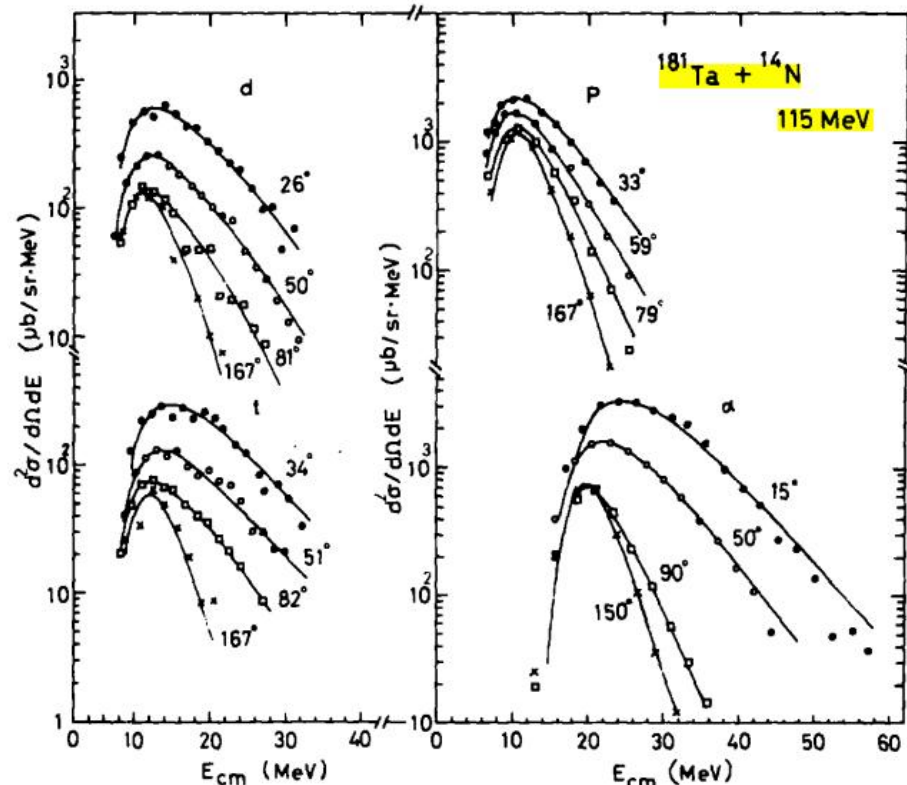


2. Preequilibrium Cluster Emission

arXiv:2505.02337



□ Results: $^{14}\text{N}+^{159}\text{Tb}$, ^{169}Tm , ^{181}Ta



3. Formation Cross Sections of SHN

□ The Dinuclear System Model (DNS)

➤ Production cross-section of SHN:

$$\sigma_{ER}(E_{c.m.}) = \sum_{J=0}^{J_{max}} \sigma_{cap}(E_{c.m.}, J) P_{CN}(E_{c.m.}, J) W_{sur}(E_{c.m.}, J)$$

$$\begin{aligned} & \frac{dP(Z_1, N_1, E_1, \beta_1, B, t)}{dt} \\ &= \sum_{Z'_1=Z_1 \pm 1} W_{Z_1, N_1, \beta_1; Z'_1, N'_1, \beta'_1}^p(t) \times [d_{Z_1, N_1} P(Z'_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z'_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \\ &+ \sum_{N'_1=N_1 \pm 1} W_{Z_1, N_1, \beta_1; Z_1, N'_1, \beta'_1}^n(t) \times [d_{Z_1, N_1} P(Z_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \\ &+ \sum_{Z'_1=Z_1 \pm 1, N'_1=N_1 \pm 1} W_{Z_1, N_1, \beta_1; Z'_1, N'_1, \beta'_1}^d(t) \times [d_{Z_1, N_1} P(Z'_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z'_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \\ &+ \sum_{Z'_1=Z_1 \pm 1, N'_1=N_1 \pm 2} W_{Z_1, N_1, \beta_1; Z'_1, N'_1, \beta'_1}^t(t) \times [d_{Z_1, N_1} P(Z'_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z'_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \\ &+ \sum_{Z'_1=Z_1 \pm 2, N'_1=N_1 \pm 1} W_{Z_1, N_1, \beta_1; Z'_1, N'_1, \beta'_1}^{3He}(t) \times [d_{Z_1, N_1} P(Z'_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z'_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \\ &+ \sum_{Z'_1=Z_1 \pm 2, N'_1=N_1 \pm 2} W_{Z_1, N_1, \beta_1; Z'_1, N'_1, \beta'_1}^{\alpha}(t) \times [d_{Z_1, N_1} P(Z'_1, N'_1, E'_1, \beta'_1, B, t) - d_{Z'_1, N'_1} P(Z_1, N_1, E_1, \beta_1, B, t)] \end{aligned}$$

$$W_{Z_1, N_1; Z'_1, N'_1}^{\nu} = G_{\nu} \frac{\tau_{mem}(Z_1, N_1, E_1; Z'_1, N'_1, E'_1)}{d_{Z_1, N_1} d_{Z'_1, N'_1} \hbar^2} \times \sum_{ii'} |\langle Z'_1, N'_1, E'_1, i' | V | Z_1, N_1, E_1, i \rangle|^2.$$

➤ σ_{cap} : the probability that the projectile nucleus and the target nucleus overcome the Coulomb barrier and form a DNS.

➤ V.I. Zagrebaev, Phys. Rev. C 64 (2001) 034606.

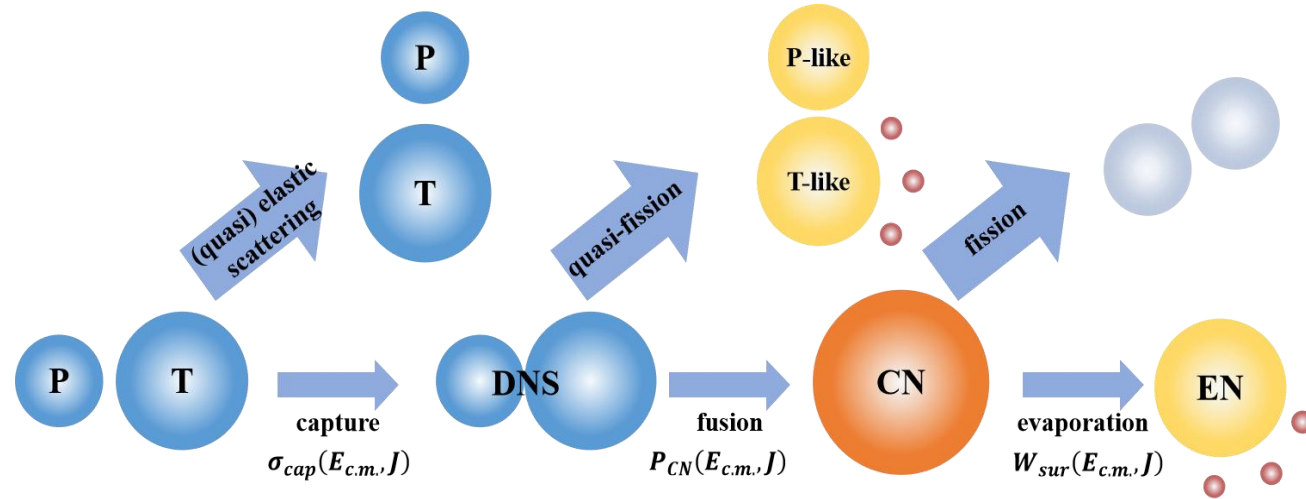
➤ P_{CN} : the formation probability of compound nuclei.

➤ N. Antonenko et al, Phys. Lett. B 319 (1993) 425.

➤ N.V. Antonenko et al, Phys. Rev. C 51 (1995) 2635.

➤ W_{sur} : the probability of the highly excited compound nuclei surviving by evaporating light particles.

➤ A.S. Zubov et al, Phys. Rev. C 65 (2002) 024308.



Schematic diagram of fusion evaporation reactions in the DNS model

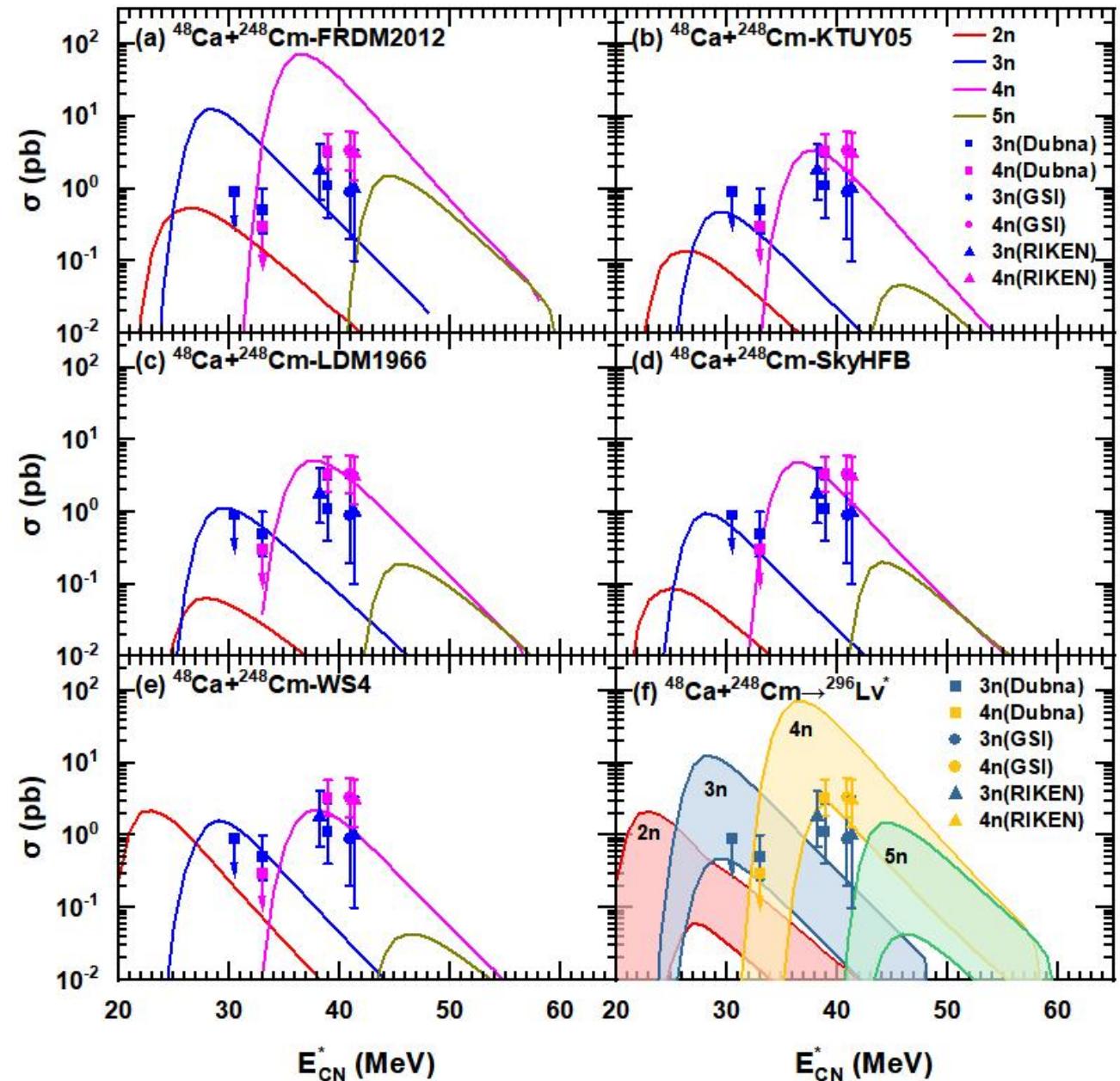
➤ 冯兆庆. 近垒重离子熔合反应和超重核合成机制研究[D]. 甘肃:中国科学院近代物理研究所, 2007.

3. Formation Cross Sections of SHN



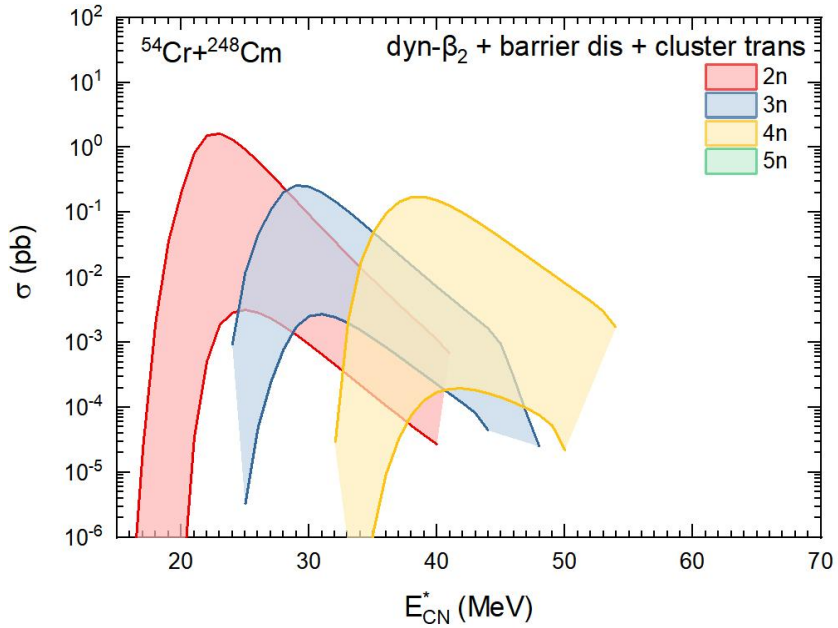
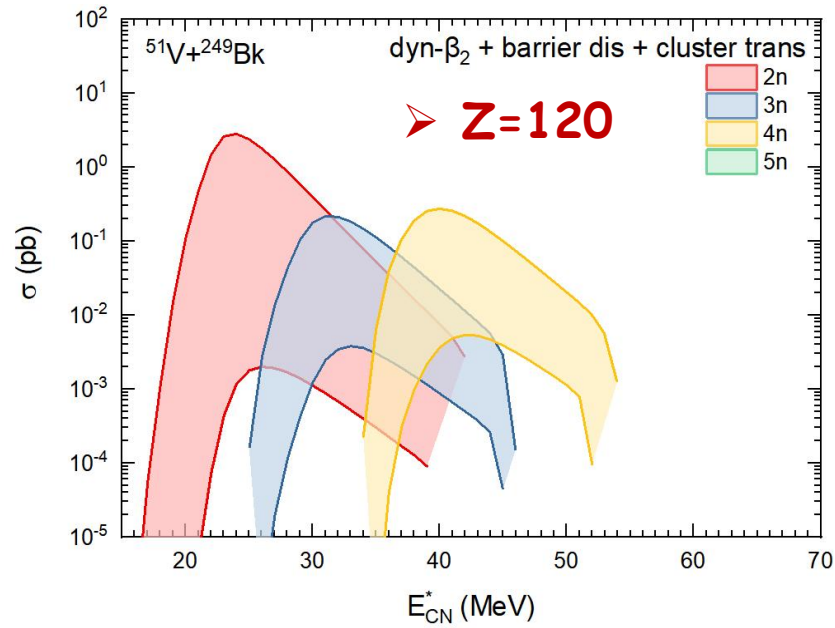
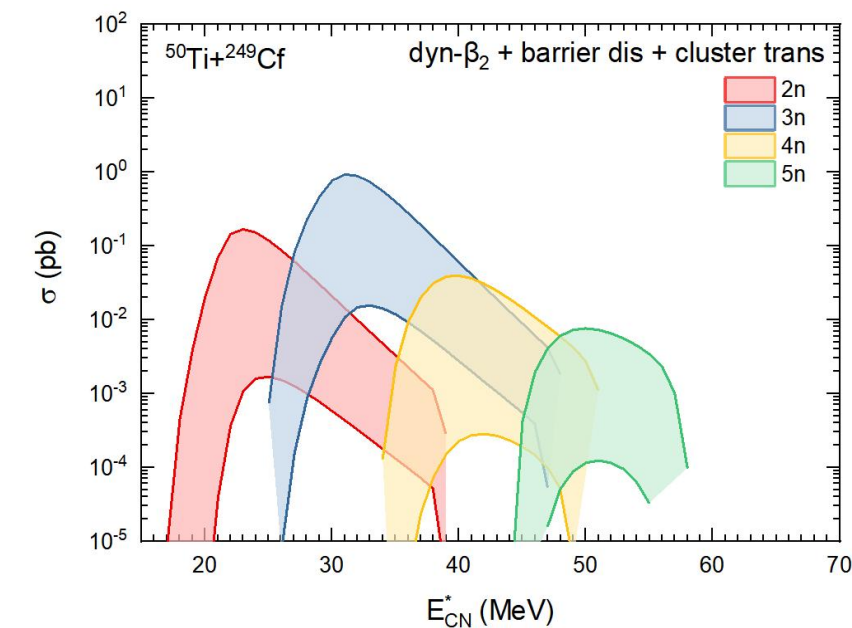
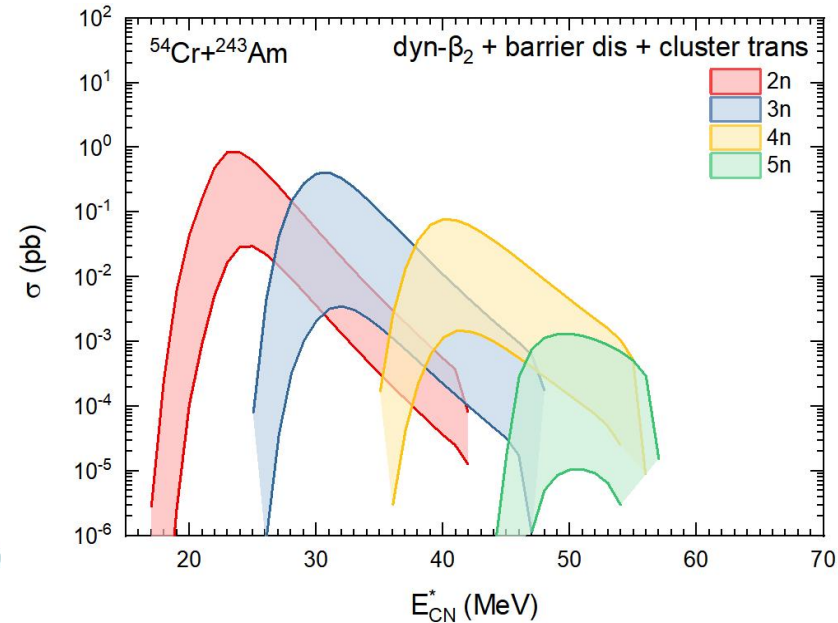
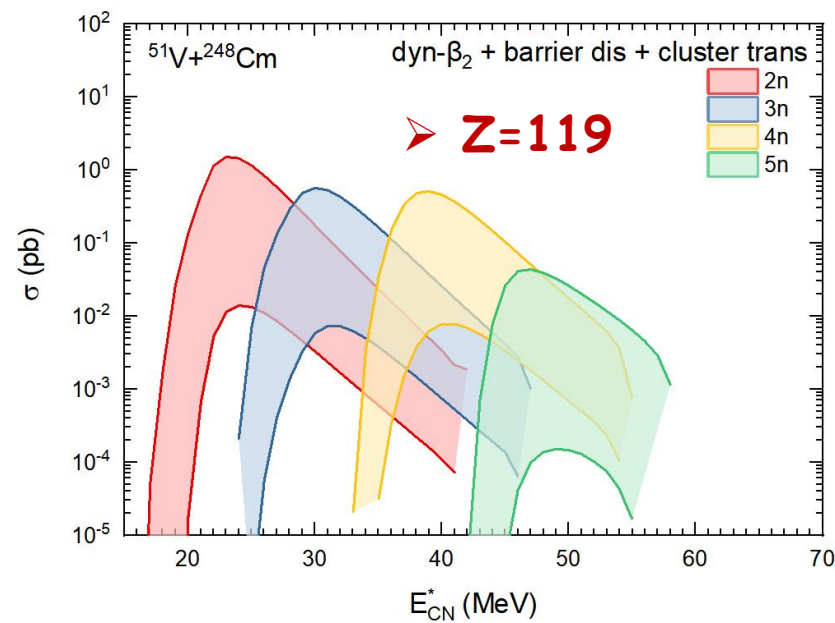
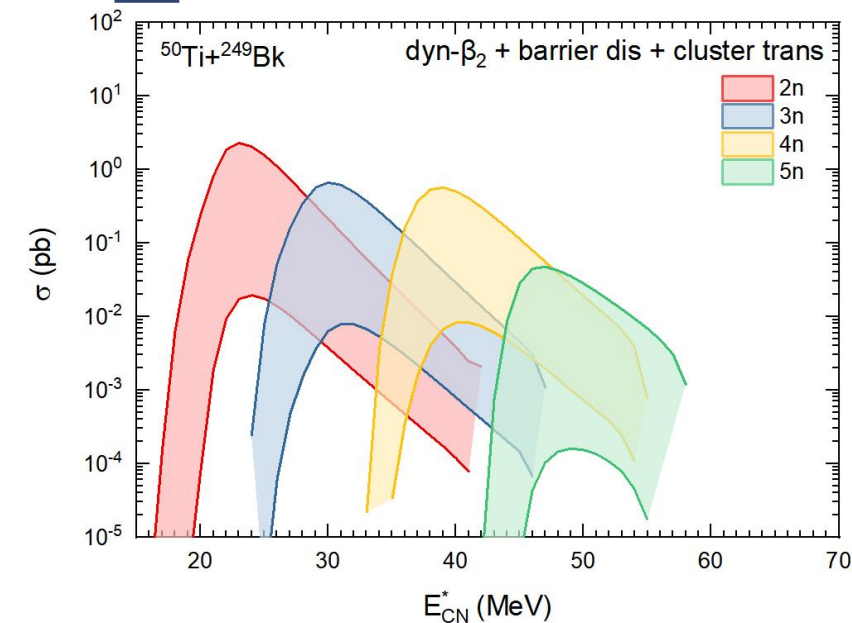
□ Different mass tables

- **FRDM2012:**
the finite-range droplet model
➤ P. Möller et al, Atomic Data and Nuclear Data Tables 109–110 (2016) 1–204
- **KTUY05:**
the Koura-Tachibana-Uno-Yamada
➤ Hiroyuki Koura et al, Progress of Theoretical Physics 113 (2005) 305–325
- **LDM1966:**
the little-droplet model
➤ William D. Myers et al, Nuclear Physics 81 (1966) 1–60
- **SkyHFB:**
the Hartree-Fock-Bogoliubov method
➤ Y. El Bassem et al, Nuclear Physics A 957 (2017) 22–32
- **WS4:**
the Weizsäcker-Skyrme mass model
➤ Ning Wang et al, Physics Letters B 734 (2014) 215–219





3. Formation Cross Sections of SHN



THANK YOU !

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