

Resumming transverse observables for NNLO+PS matching in GENEVA

Alessandro Gavardi, **Rebecca von Kuk**, Matthew Lim, [2505.14773](#)

Resumming transverse observables for NNLO+PS matching in GENEVA for the $b\bar{b}H$ channel

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1. Motivation

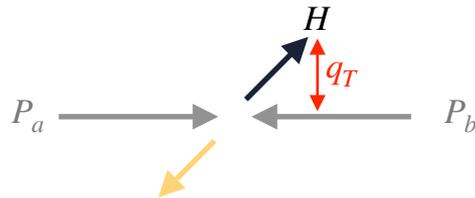
2. Resolution variables q_T and $\mathcal{T}_1^{p_T}$

3. Resummation of relevant observables

- q_T resummation in SCET
- Resummation of generalized one-jettiness $\mathcal{T}_1^{p_T}$

4. Results for $b\bar{b}H$

5. Outlook and Summary



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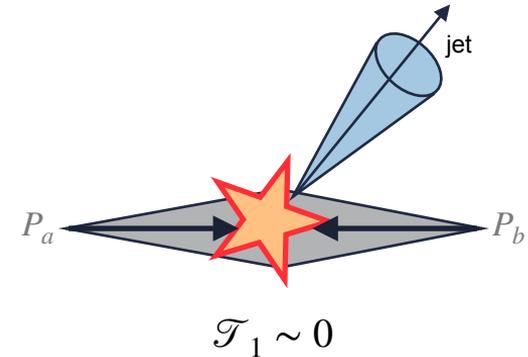
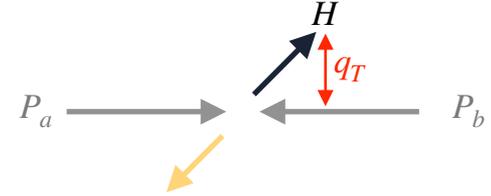
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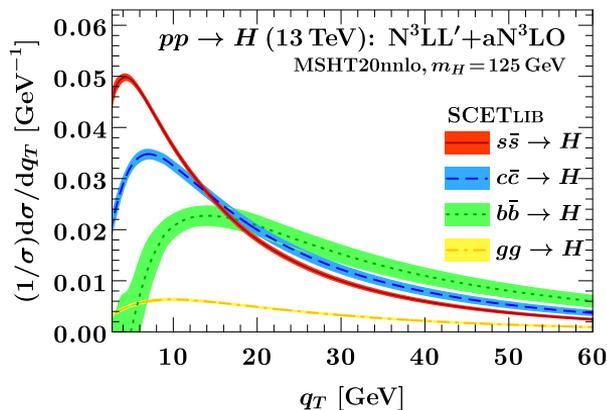


Introduction

- LHC precision demands theoretical tools at highest accuracy
- Event generators provide important bridge between theory and experiment
 - Simulate process from hard interaction to final states that can be detected at colliders
- State-of-art for predictions matched to parton showers: NNLO+PS
 - **GENEVA** [Alioli et al., '13, '15]
 - **MiNNLO_{PS}** [Monni, Nason, Re, Wiesemann, Zanderighi, '20]
- Study different choices for jet resolution variables

Quark initiated Higgs production

- Higgs q_T spectrum allows to access Yukawa coupling from Higgs production
- Allows initial state discrimination [Ebert et al. '16, Bishara et al. '16]
 - q_T spectra of ggH and qqH have different shapes
- High precision prediction available [Cal, RvK, Lim, Tackmann '23]
- NNLO+PS prediction in MiNNLO_{PS} [Biello, Sankar, Wiesemann, Zanderighi '24]
- **Now:** implementation in GENEVA with transverse observables



Resolution variables

Calculations beyond leading order

- IR divergences cancel between matrix elements with different numbers of final states

$$\sigma^{\text{NLO}} = \int d\Phi_N [B_N(\Phi_N) + V_N(\Phi_N)] + \int d\Phi_{N+1} [B_{N+1}(\Phi_{N+1})]$$

$$\frac{A}{\epsilon^2} + \frac{B}{\epsilon} + C$$

$$-\frac{A}{\epsilon^2} - \frac{B}{\epsilon} + D$$

- MC: cancellation achieved numerically by defining appropriate subtraction terms
- In GENEVA: convert IR-divergent final state with $N + 1$ partons into IR-finite final state with N partons

Resolution variables

- Resolution variables r : divide phase space into bins: $r < r^{\text{cut}}$
 - Examples: N-jettiness (\mathcal{T}_N) [Stewart, Tackmann, Waalewijn '10], q_T , ...
 - Need to be resummed to high orders
- Predictions for exclusive observables depend on choice of resolution variable
 - Compare prediction with different resolution variables to gauge uncertainties
- Higher-accuracy parton showers are becoming available [Dasgupta, Dreyer, Hamilton, Monni, Salam, Sodes '20, Herren, Höche, Krauss, Reichelt, Schoenherr '23, van Beekveld et al. '24, Höche, Krauss, Reichelt '24]
 - Explore novel variable choices that preserve shower accuracy



Resolution variables

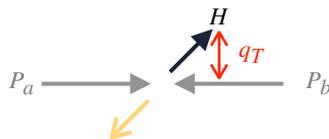
Original resolution variables in GENEVA:

\mathcal{T}_0 and \mathcal{T}_1

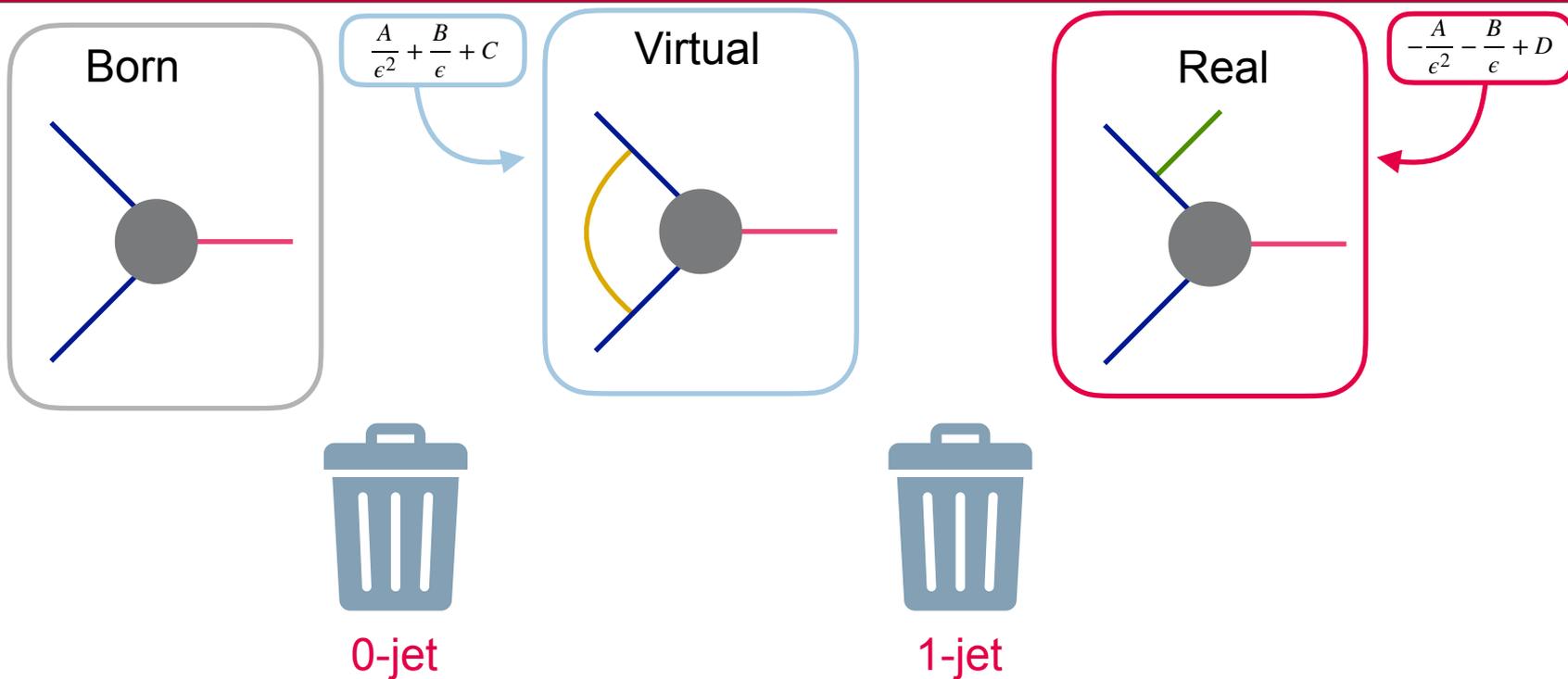
- Simple factorization structure
- **+** Facilitates resummation
- Usually uses invariant mass-like measure
- **-** complicates matching to PS

Here: transverse observables to separate bins

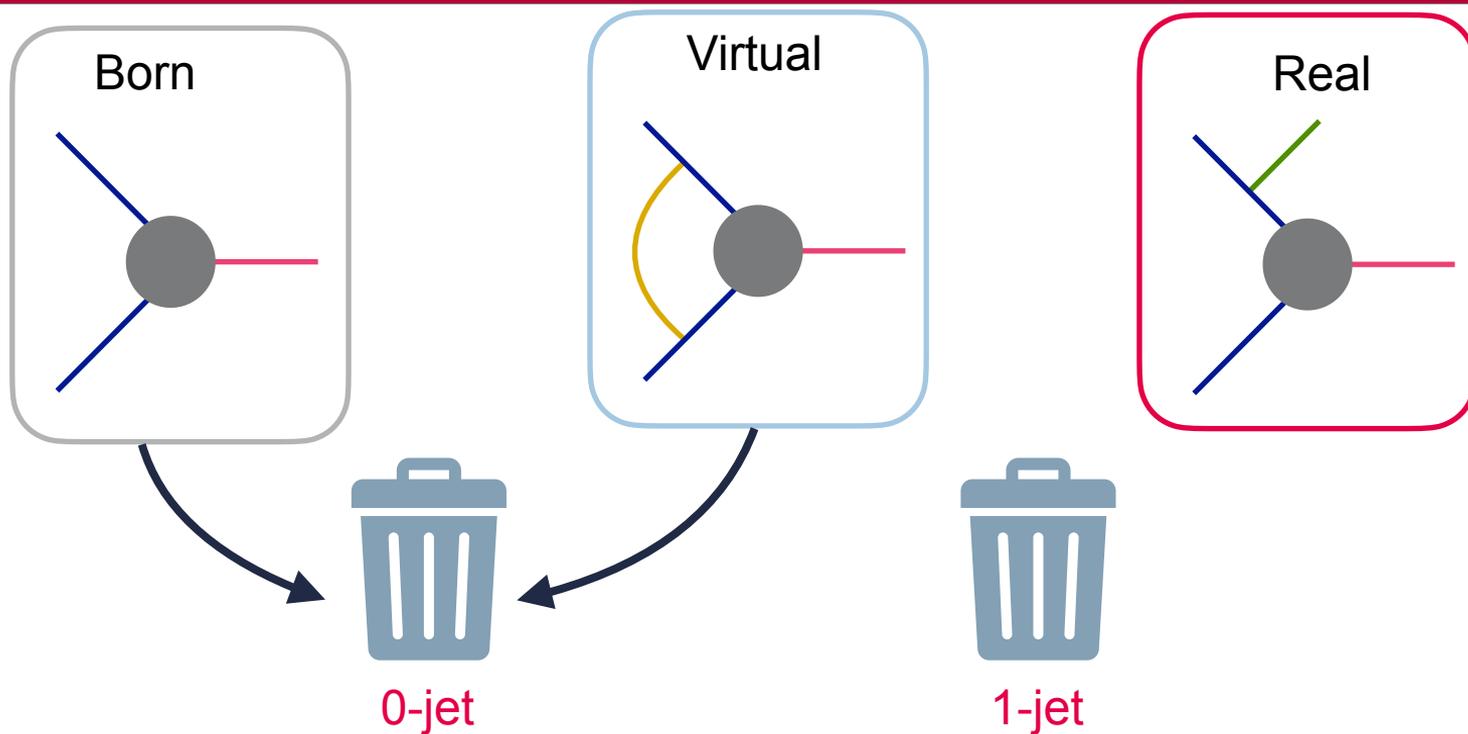
- Transverse momentum q_T to separate 0/1-jet bin
- \mathcal{T}_1 with transverse momentum-like measure
- **+** Facilitates matching to PS which tend to be ordered in transverse momentum



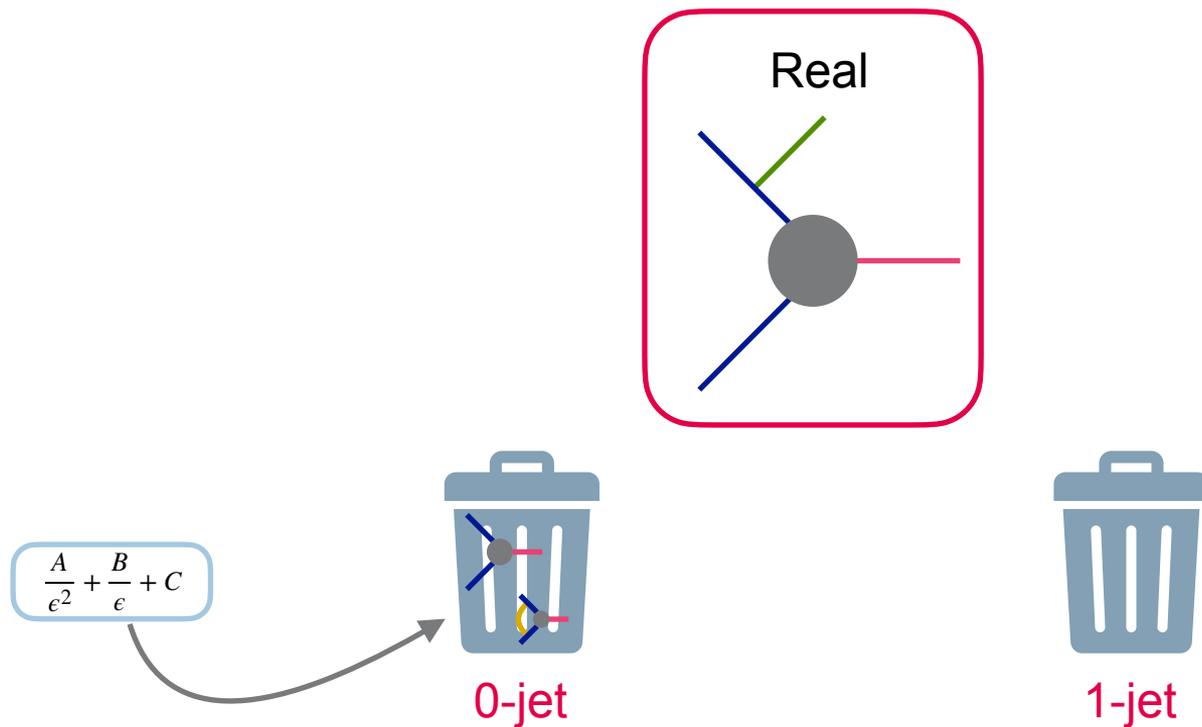
Defining IR safe events



Defining IR safe events

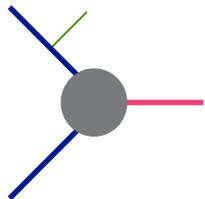


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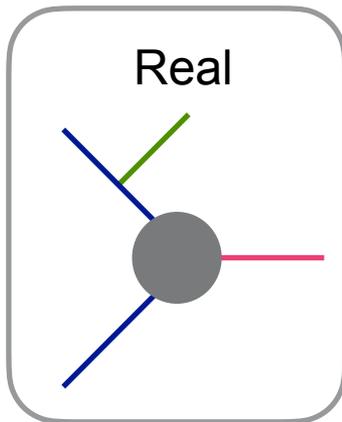
Defining IR safe events

$$-\frac{A}{\epsilon^2} - \frac{B}{\epsilon} + D$$

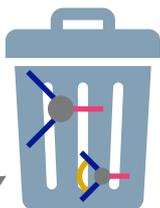


soft/collinear

$$q_T < q_T^{\text{cut}}$$



$$\frac{A}{\epsilon^2} + \frac{B}{\epsilon} + C$$

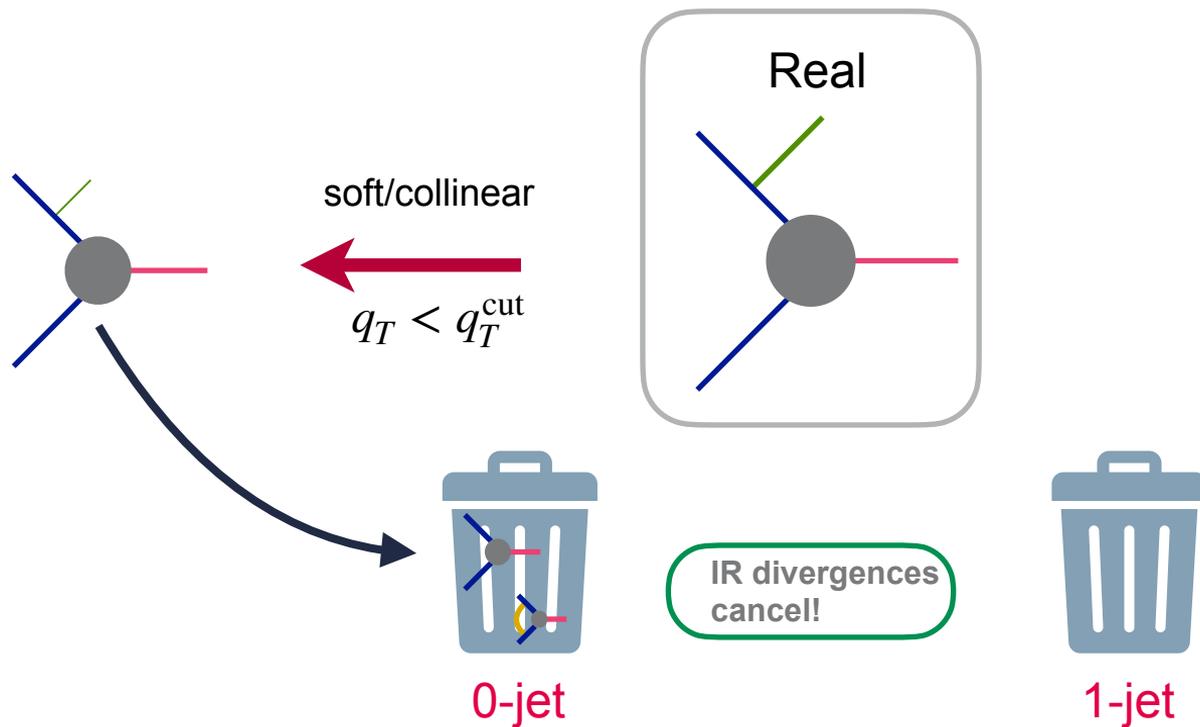


0-jet

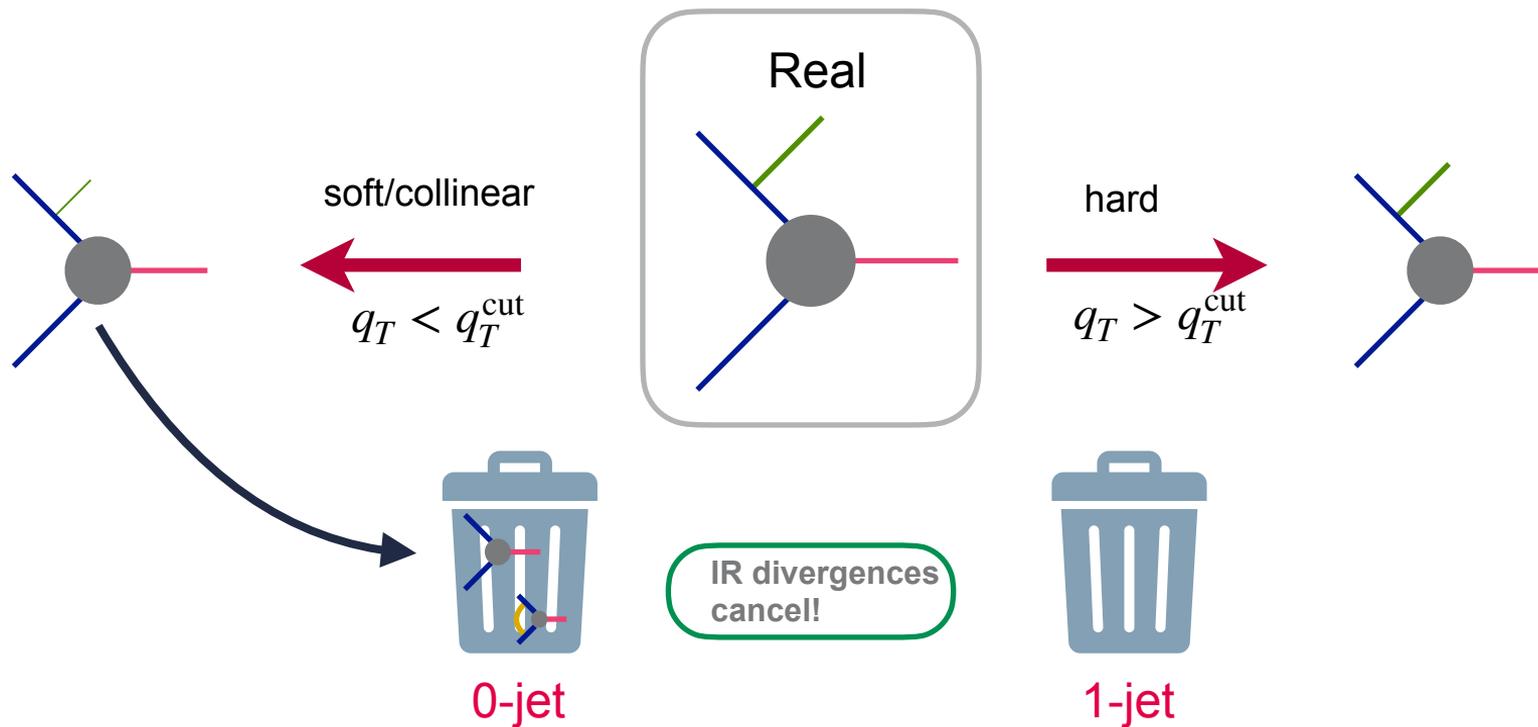


1-jet

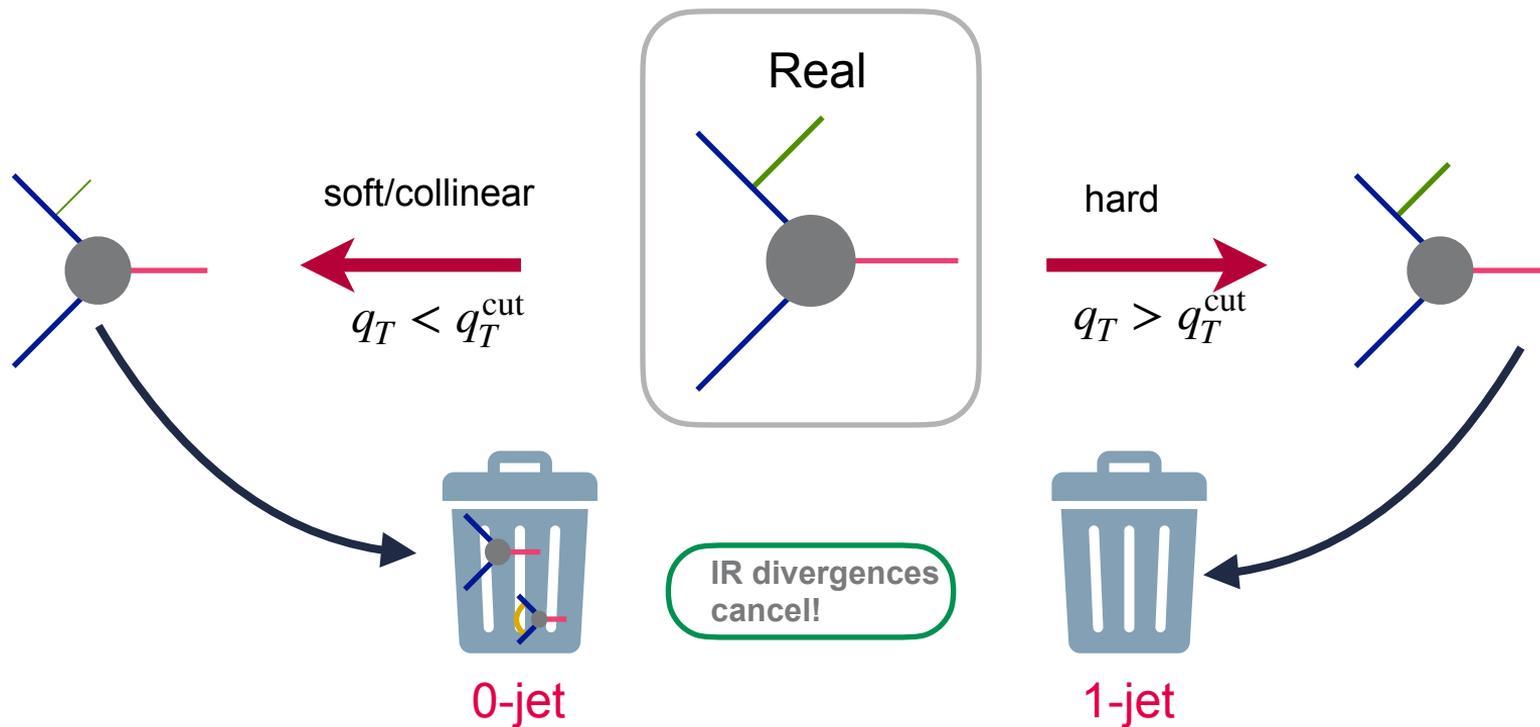
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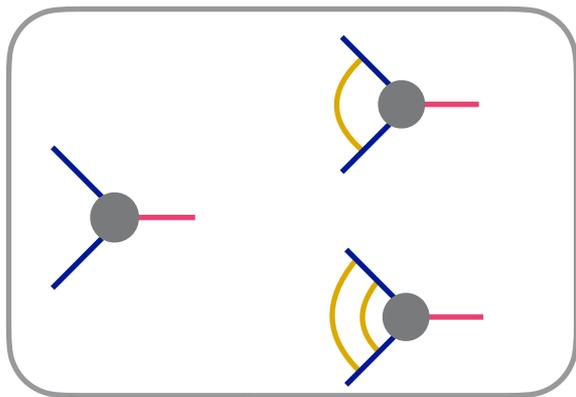
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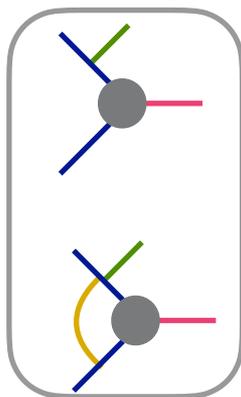
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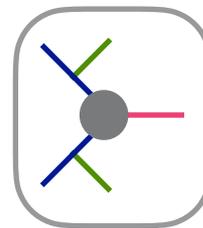
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0-jet

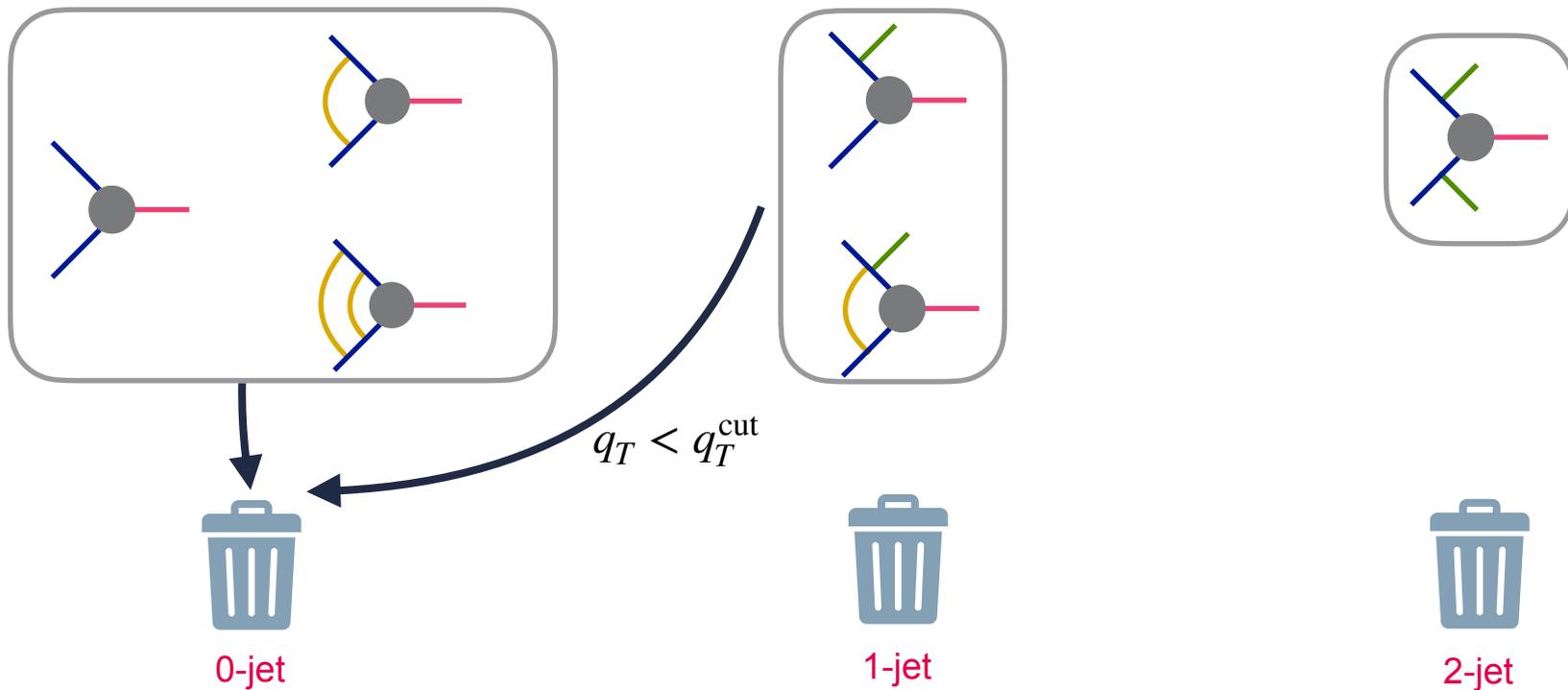


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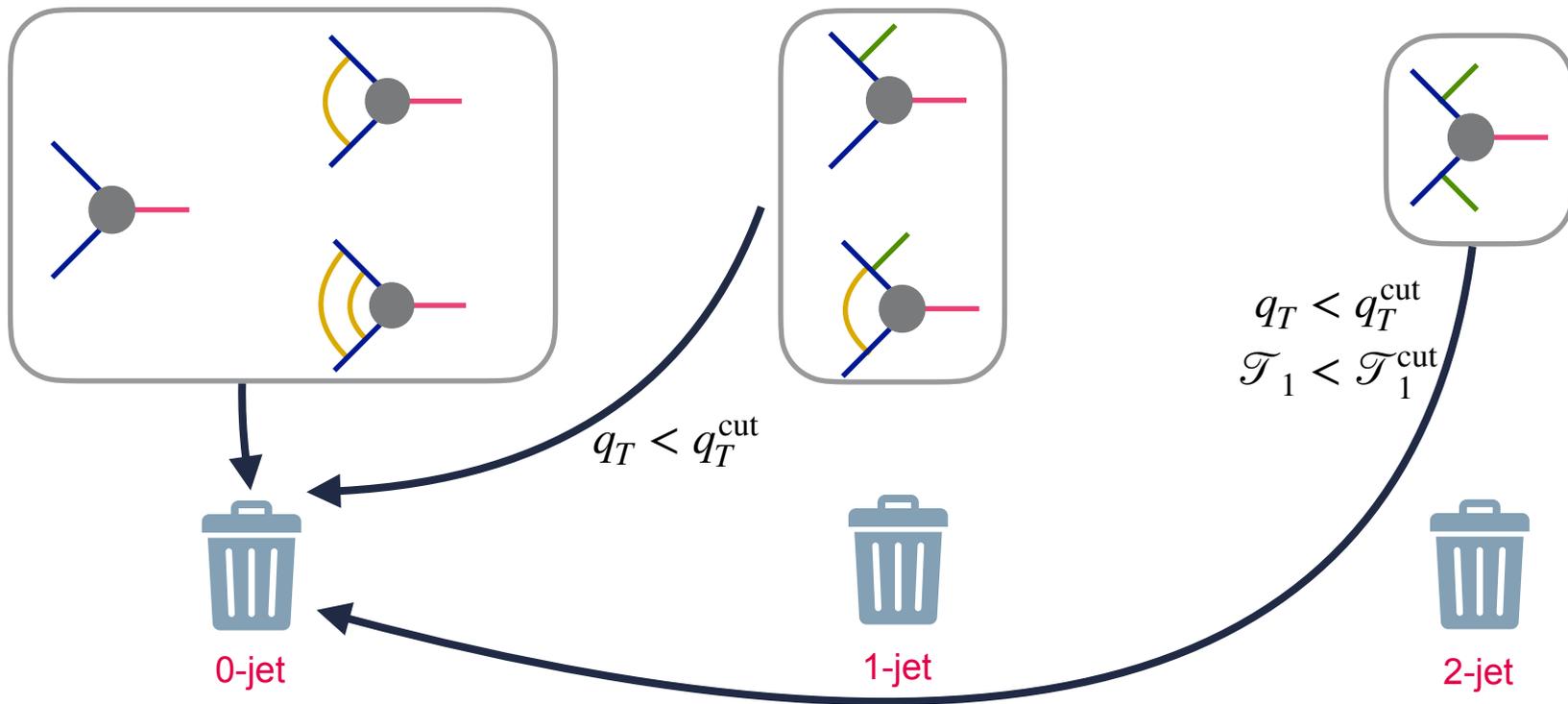


2-jet

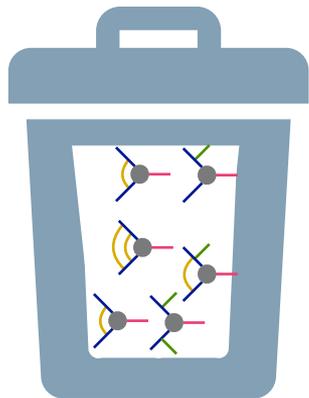
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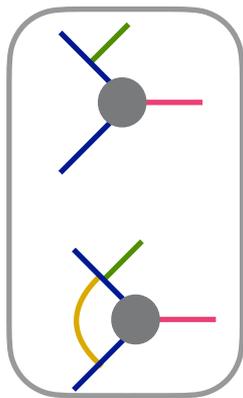
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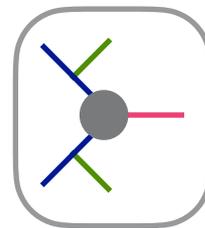
0-jet



$q_T > q_T^{\text{cut}}$

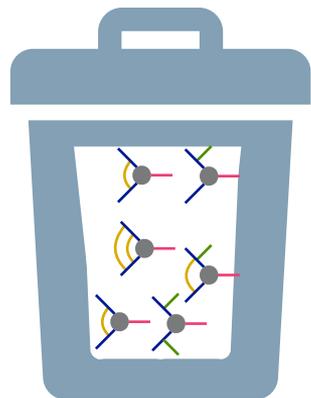


1-jet

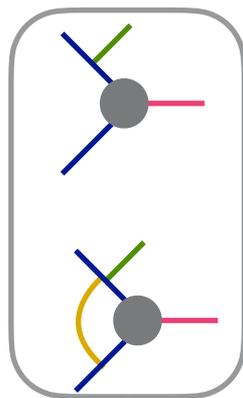


2-jet

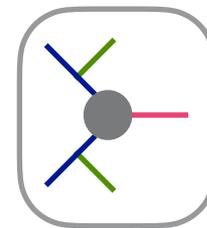
Defining IR safe events



0-jet



1-jet

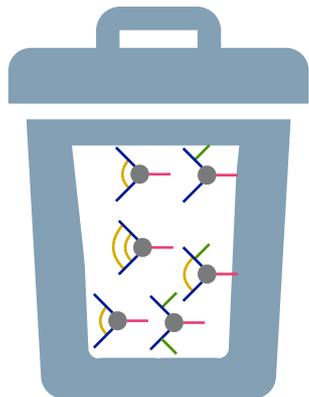


2-jet

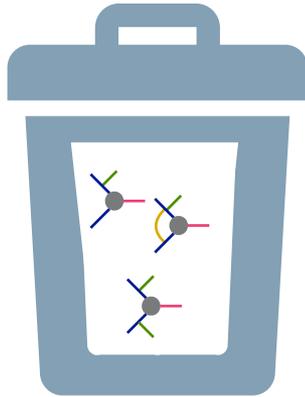
$$q_T > q_T^{\text{cut}}$$
$$\mathcal{T}_1 < \mathcal{T}_1^{\text{cut}}$$

Defining IR safe events

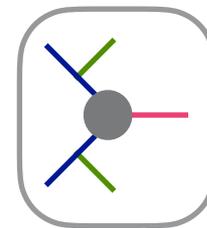
$$q_T > q_T^{\text{cut}}, \mathcal{T}_1 < \mathcal{T}_1^{\text{cut}}$$



0-jet



1-jet



$$q_T > q_T^{\text{cut}} \\ \mathcal{T}_1 > \mathcal{T}_1^{\text{cut}}$$

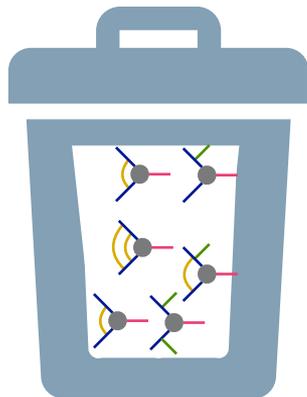


2-jet

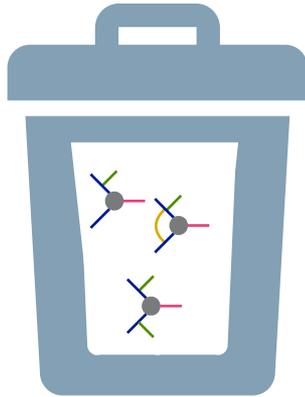
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$$q_T > q_T^{\text{cut}}, \mathcal{T}_1 < \mathcal{T}_1^{\text{cut}}$$

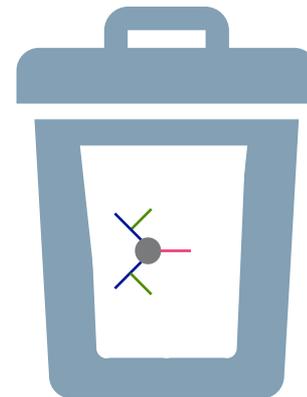
$$q_T > q_T^{\text{cut}}, \mathcal{T}_1 > \mathcal{T}_1^{\text{cut}}$$



0-jet



1-jet



2-jet



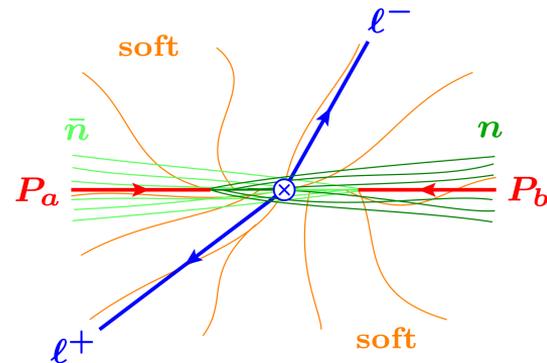
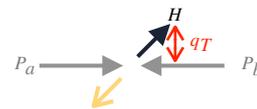
Resumming q_T and \mathcal{T}_1

q_T factorization

- Factorization in soft-collinear effective theory (SCET):

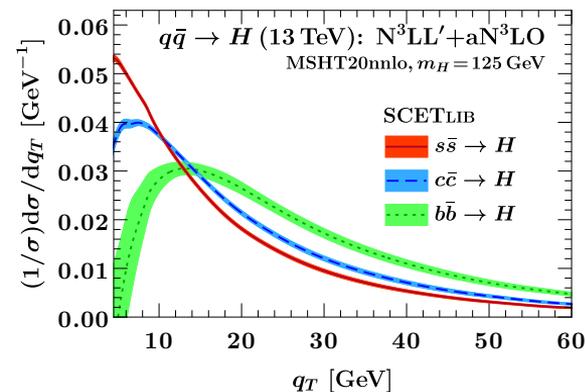
$$\frac{d\sigma}{dq_T} = H(Q, \mu) \times B(q_T, \mu) \otimes B(q_T, \mu) \otimes S(q_T, \mu) \left[1 + \mathcal{O}\left(\frac{q_T^2}{Q^2}\right) \right]$$

- Hard function: hard interaction at high scale Q
- Beam function: collinear radiation
- Soft function: soft, isotropic radiation



q_T resummation

- Solve RGEs for $H(Q, \mu)$, $B(q_T, \mu)$ and $S(q_T, \mu)$ to resum large logs of q_T/Q
- In practice: use SCETlib [Ebert, Michel, Tackmann '18] to perform resummation
- **New:** directly interface SCETlib to GENEVA
- Resummation for $q\bar{q}H$ available up to N³LL`
 - $H(Q, \mu)$ [Gehrmann, Kara '14; Ebert, Michel, Tackmann '17]
 - $B(q_T, \mu)$ [Luo, Yang, Zhu, Zhu '19; Ebert, Mistlberger, Vita '20]
 - $S(q_T, \mu)$ [Li, Zhu, Neill '16; Li, Zhu, '16]
 - 4-loop cusp and 3-loop non-cusp [Henn, Korchemsky, Mistlberger '20; v. Manteuffel, Panzer, Schabinger '20] [Li, Zhu, '16; Vladimirov '16]

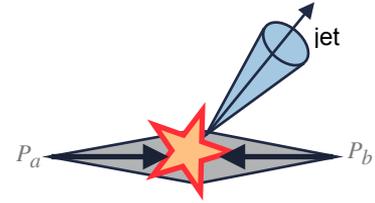


[Cal, RvK, Lim, Tackmann '23]

- Standard definition of jettiness [Stewart, Tackmann, Waalewijn, '10]

$$\mathcal{T}_1 = \sum_i \min\{d_1(p_i), d_a(p_i), d_b(p_i)\}$$

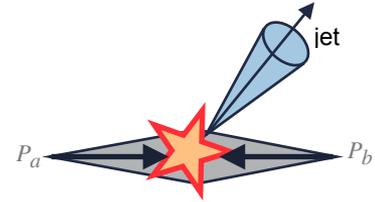
colored final states p_i , $m \in \{1, \dots, N + 2\} = \{1, 2, 3\}$, distance measure to region m $d_m(p_i)$



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colored final states p_i , $m \in \{1, \dots, N+2\} = \{1, 2, 3\}$, distance measure to region m $d_m(p_i)$



- Generalized definition: measure $d_m(p_i)$ doesn't need to be identical to observable f_m

[Stewart, Tackmann, Waalewijn, '10]

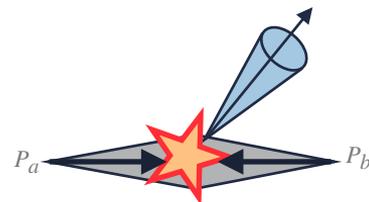
$$\mathcal{T}_1 = \sum_m \sum_{i \in m} f_m(\eta_i, \phi_i) p_{T,i}$$

observable f_m , pseudo rapidity η_i , azimuth ϕ_i and transverse momentum $p_{T,i}$

Generalized \mathcal{T}_1 : p_T -like measure

- Define generalized \mathcal{T}_1 with transverse momentum-like measure where $\mathcal{T}_1 \sim p_{T,i} \sim Q\lambda$

$$\mathcal{T}_1^{p_T} = \sum_i \begin{cases} p_{T,i}, & \text{for } d_B(p_i) < d_J(p_i), \\ p_{T,i} f_J(\eta_i, \phi_i), & \text{for } d_J(p_i) < d_B(p_i), \end{cases}$$



- Beam region: p_T is sufficient as observable definition
- Jet region: emission collinear to jet axis scale \Rightarrow need angular projection

\Rightarrow Use boost-invariant generalization of broadening $f_J(\eta_i, \phi_i) = \sqrt{2 \cosh(\Delta\eta_{iJ}) - 2 \cos(\Delta\phi_{iJ})}$

[Banfi, Salam, Zanderighi '04, Bertolini, Kolodrubetz, Neill, Pietrulewicz, Stewart, Tackmann, Waalewijn '17]

$$\begin{aligned} \Delta\eta_{iJ} &= \eta_i - \eta_J \\ \Delta\phi_{iJ} &= \phi_i - \phi_J \end{aligned}$$

- Choice of observables correspond to SCET_{II} setup
- Factorization theorem [Bertolini, Kolodrubtez, Neill, Pietrulewicz, Stewart, Tackmann, Waalewijn '17]

$$\frac{d\sigma}{d\mathcal{T}_1^{p_T}} = H(Q, \mu) \otimes S(\mathcal{T}_1^{p_T}, \mu) \otimes B(\mathcal{T}_1^{p_T}, \mu) \otimes B(\mathcal{T}_1^{p_T}, \mu) \otimes J(\mathcal{T}_1^{p_T}, \mu)$$

- Hard, beam and soft functions similar as for q_T factorization
- Additional **jet function** describes collinear radiation from the jet

- Solve RGEs for $H(Q, \mu)$, $B(\mathcal{T}_1^{p_T}, \mu)$, $S(\mathcal{T}_1^{p_T}, \mu)$ and $J(\mathcal{T}_1^{p_T}, \mu)$ to resum large logs
 - 1-loop soft function [Bertolini, Kolodrubtez, Neill, Pietrulewicz, Stewart, Tackmann, Waalewijn '17]
 - 1-loop beam functions [Becher, Neubert '12; Becher, Neubert, Rothen '13; Stewart, Tackmann, Walsh, Zuberi '14]
 - 1-loop recoil-free jet function [Larkoski, Neill, Thaler '14]
- Evolution equations solved in Laplace space
 - At NLL(^o): evolution equations can be solved analytically
 - Non-trivial dependence of rapidity anom. dim. on transverse momenta complicates inverse Laplace transform at higher orders

Results for $b\bar{b}H$

Fixed-order Validation

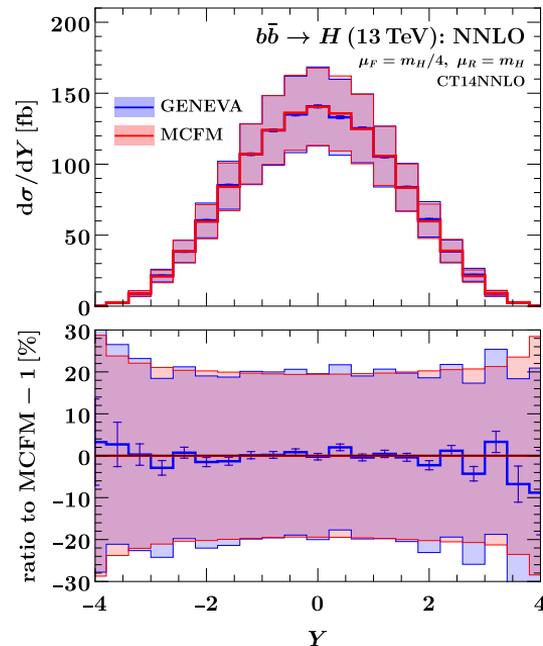
NNLO Higgs boson rapidity spectrum

- Comparing to NNLO MCFM result

[Boughezal, Campbell, Ellis, Focke, Giele, Liu, Petriello, Williams '17]

[Mondini, Williams '21]

- Good agreement with central value and 7-point scale variations

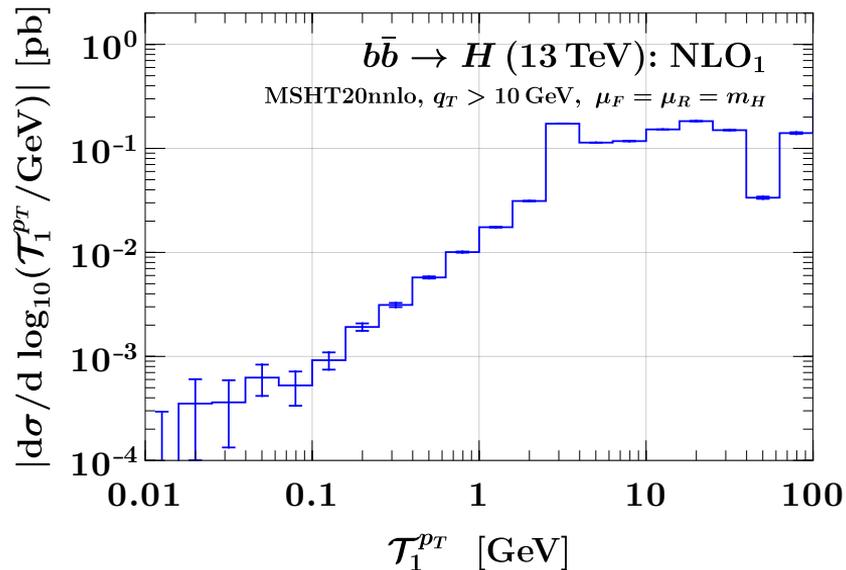


1/2-jet bin separation

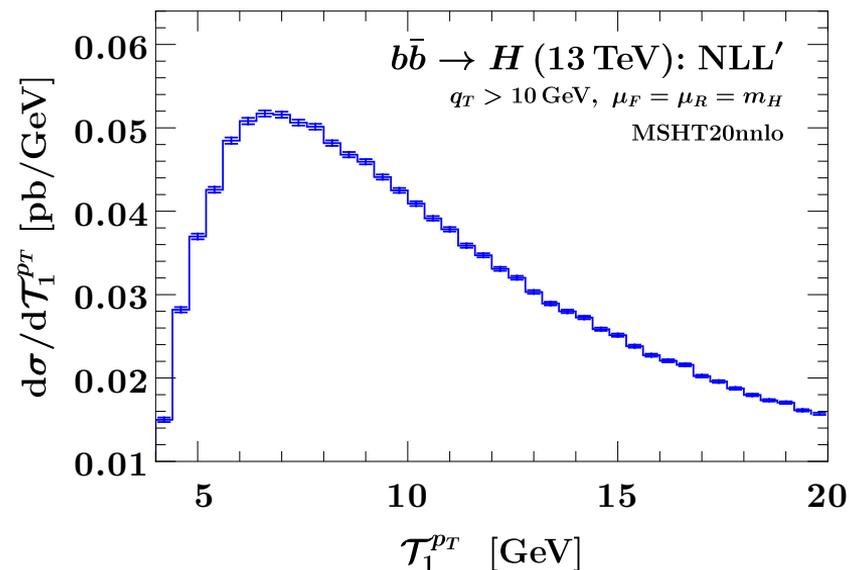
\mathcal{T}_1^{pT} spectrum

[Gavardi, RvK, Lim '25]

Non-singular



Resummed



Effect of the showering

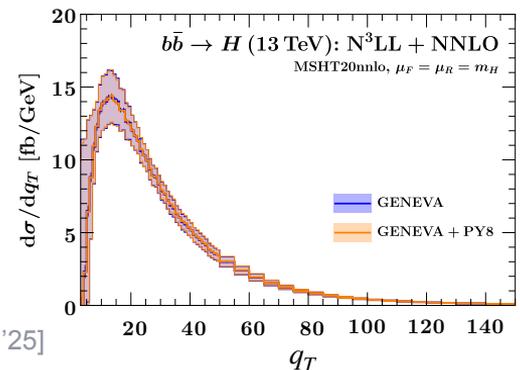
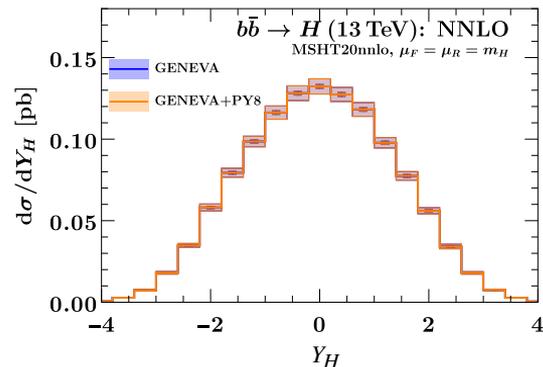
Higgs rapidity spectrum

- PS does not affect distribution
- Expected as Y_H is an inclusive quantity

Higgs transverse momentum spectrum

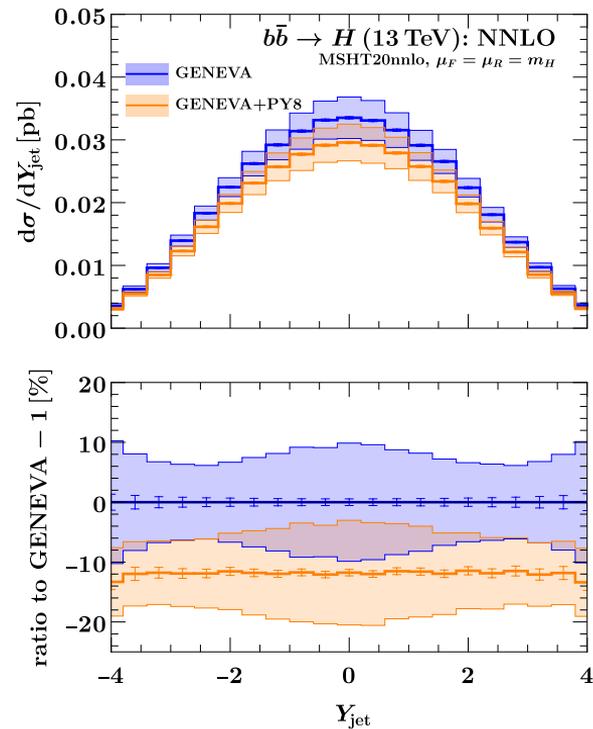
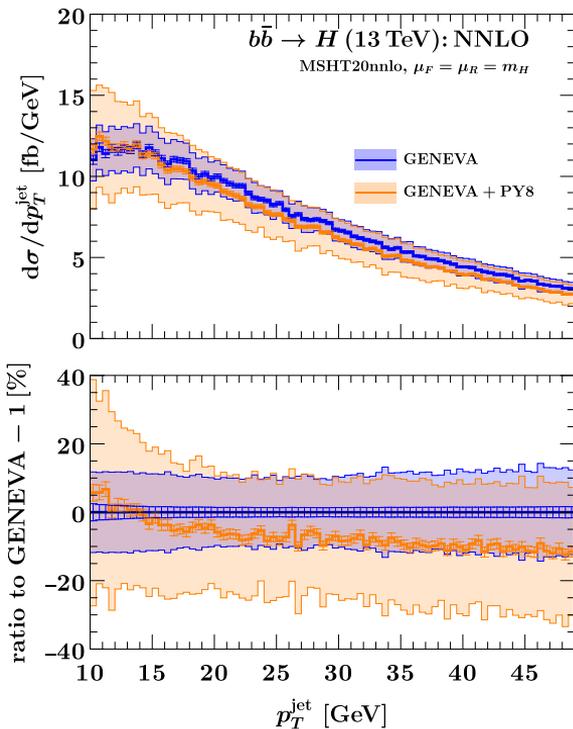
- Very good agreement between parton-level and showered results
- Can't formally claim N³LL+NNLO accuracy for showered results but distributions are indistinguishable

[Gavardi, RvK, Lim '25]



Effect of the showering

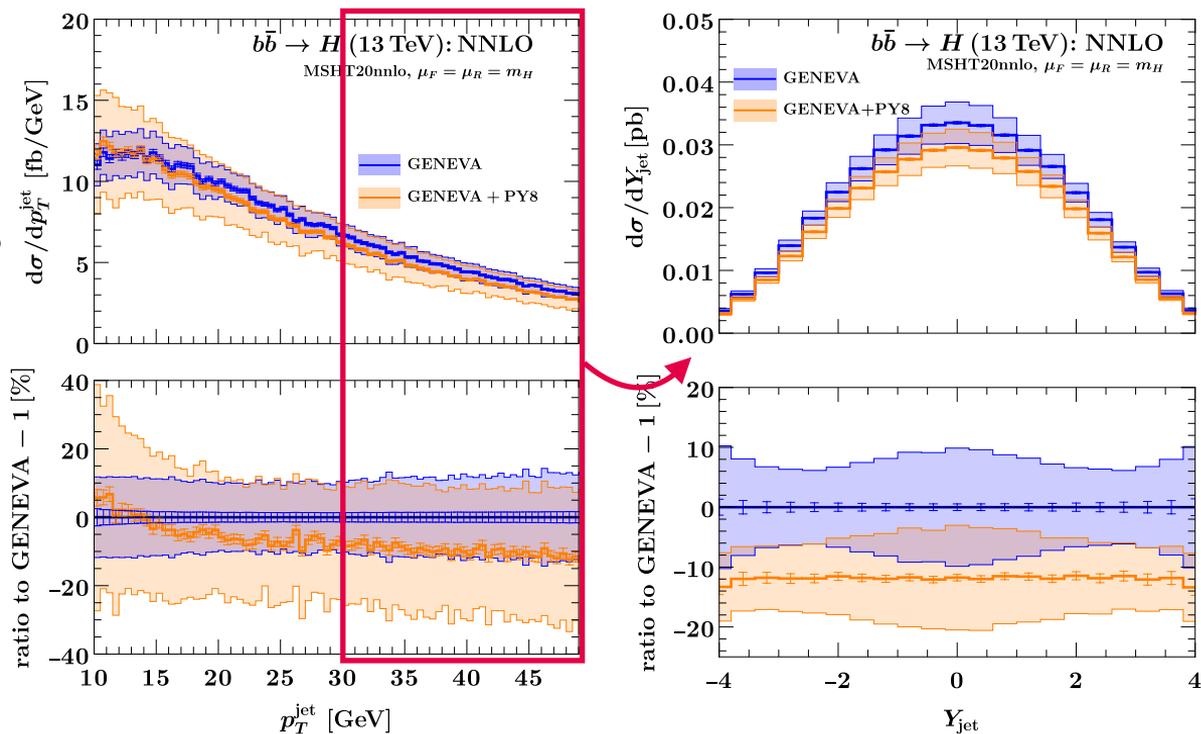
- More exclusive observables
 - p_T^{jet} : p_T of hardest jet
 - Y_{jet} : rapidity of hardest jet
- Shower effects expected



[Gavardi, RvK, Lim '25]

Effect of the showering

- More exclusive observables
 - p_T^{jet} : p_T of hardest jet
 - Y_{jet} : rapidity of hardest jet
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[Gavardi, RvK, Lim '25]

Outlook and Summary

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Outlook

- Extension to other color-singlet processes
- Achieve NNLL` resummation accuracy for $\mathcal{T}_1^{P_T}$

Summary

- New set of resolution variables in GENEVA for NNLO+PS matching
 - Facilitates matching to parton showers
- Good agreement with independent fixed-order prediction
- Resummation accuracy not affected by the shower

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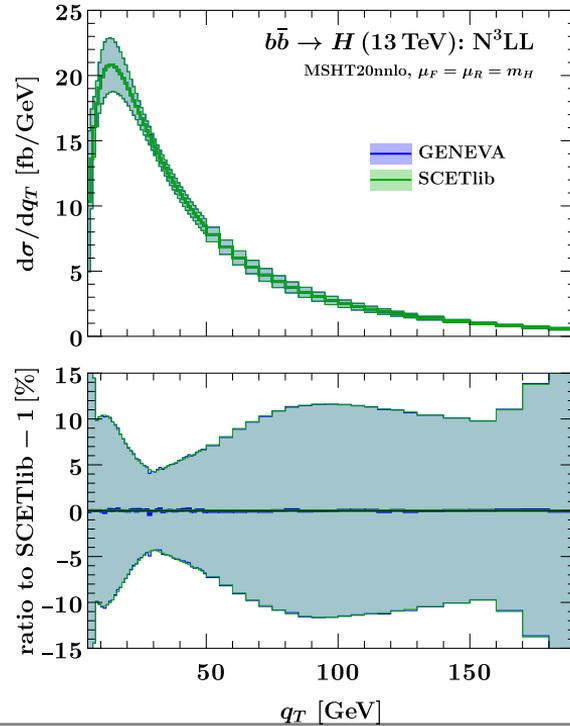
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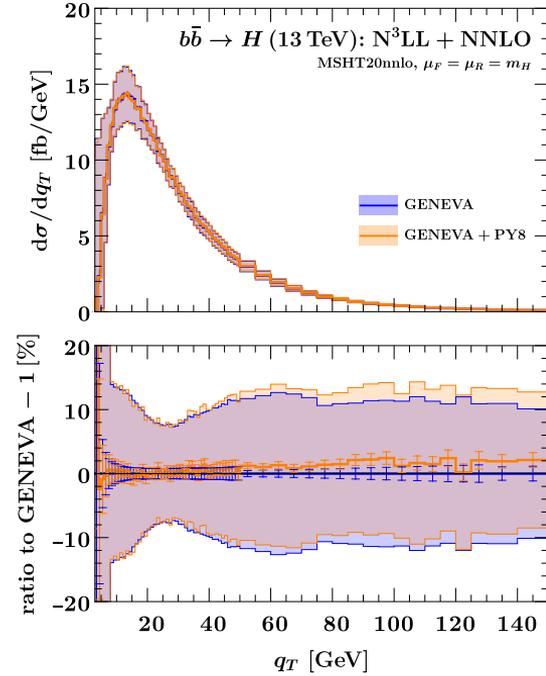
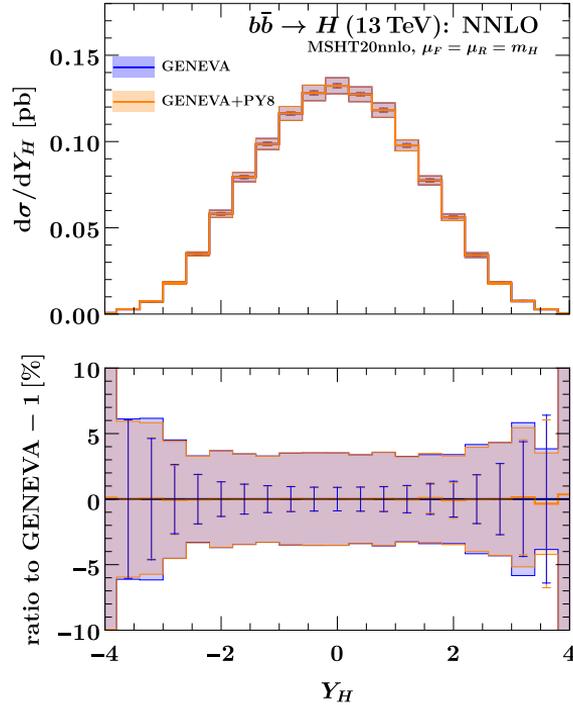
Thank you!

Back up

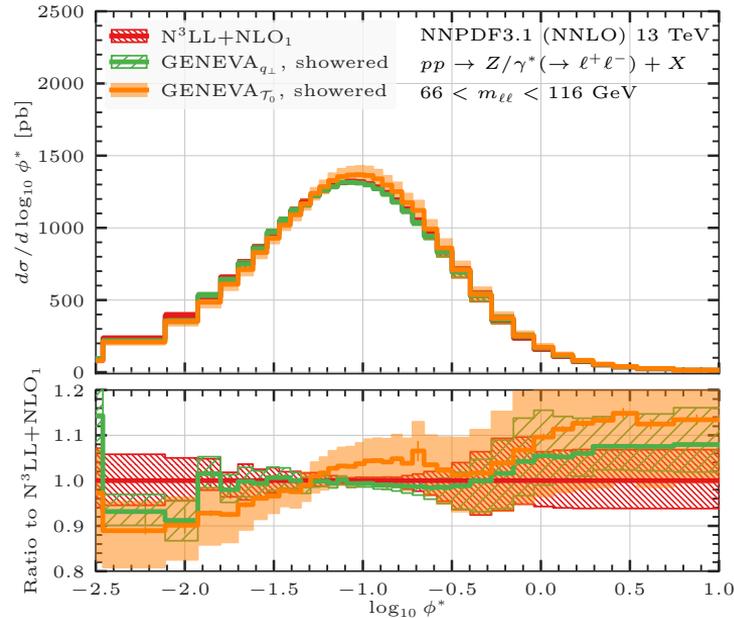
Resumed Validation



Effect of the showering



Comparison of resolution variables



[Alioli et al. '21]